

The Multi-Messenger View of AGN

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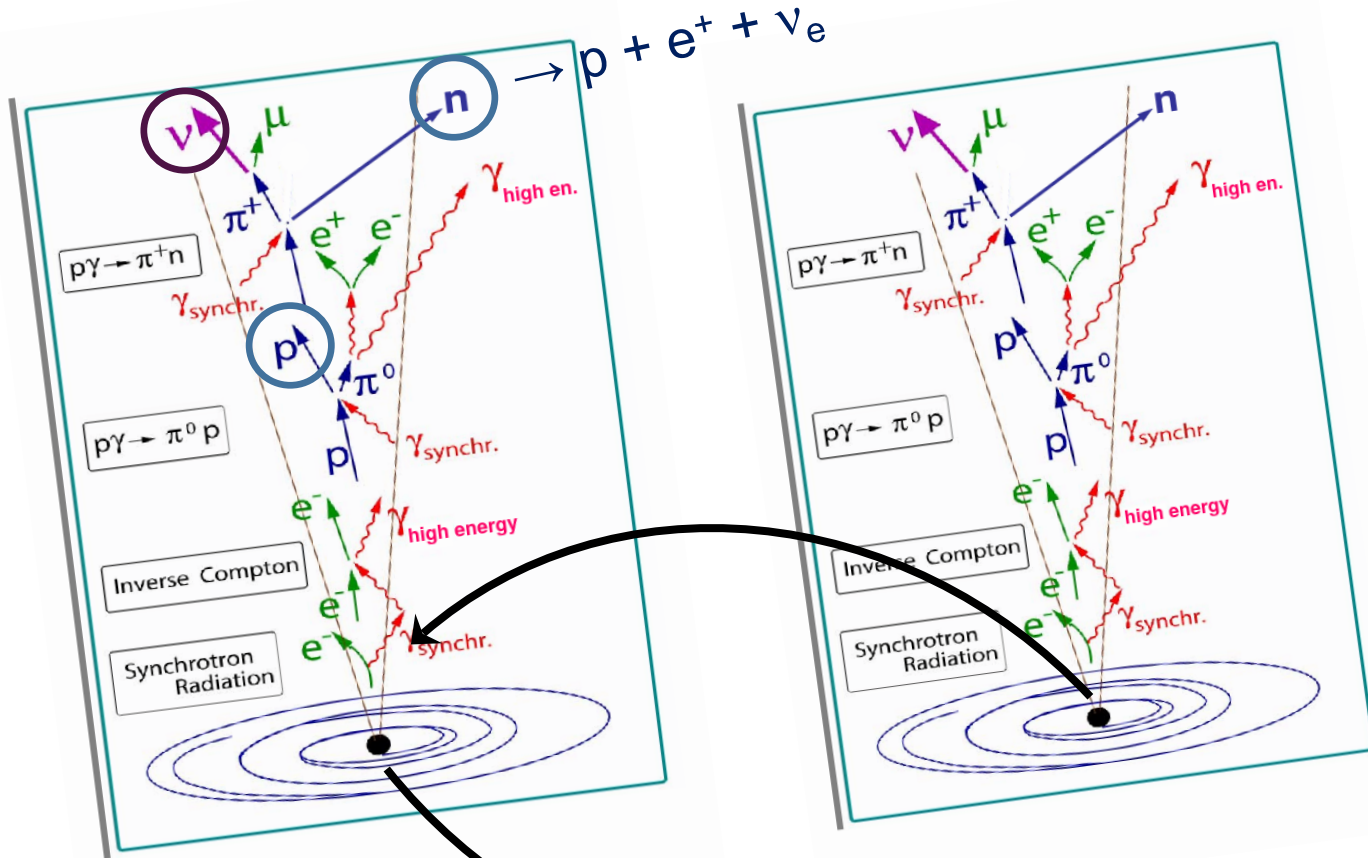


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Multi-Messenger Signals from AGN (hadronic / lepto-hadronic models)

Neutrinos

Cosmic Rays (p)

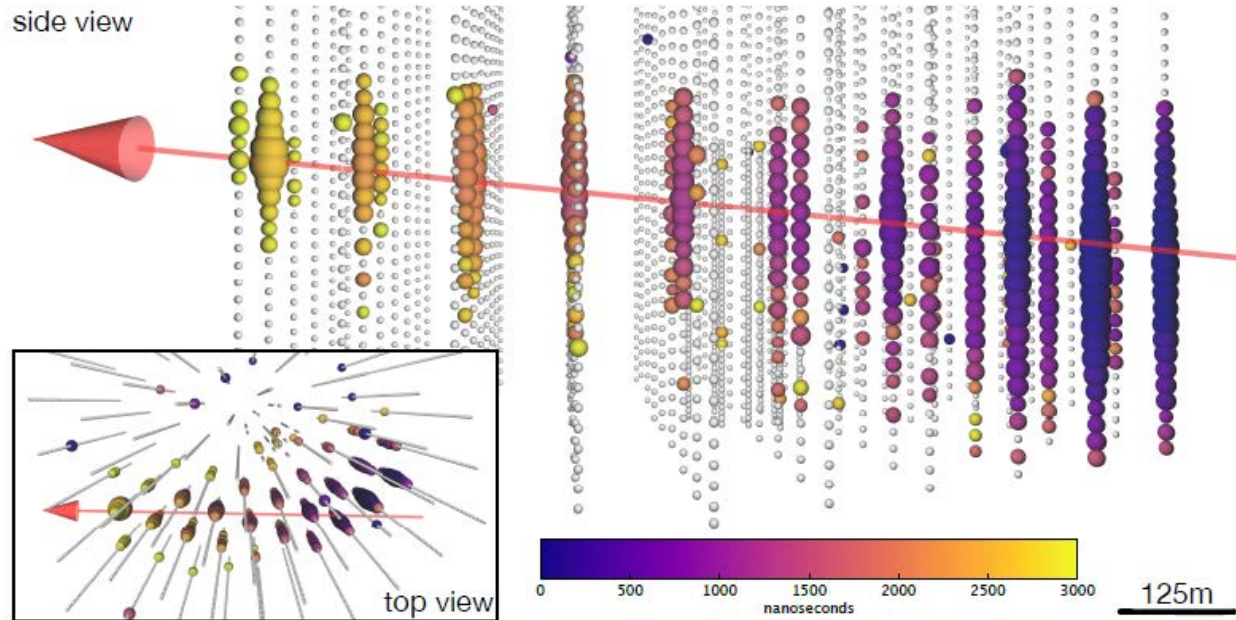


Gravitational Waves

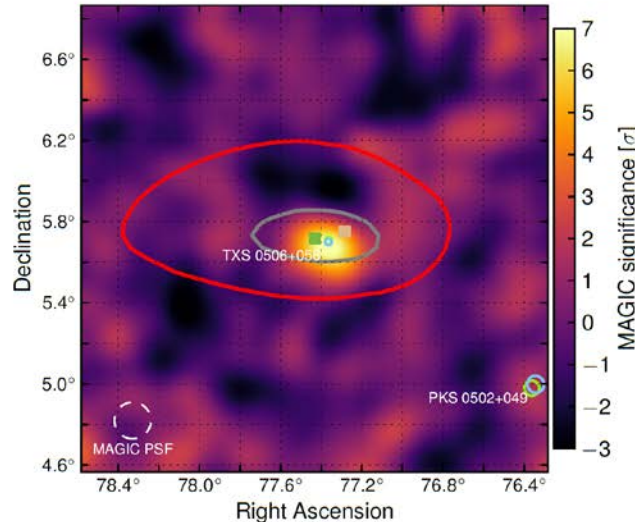
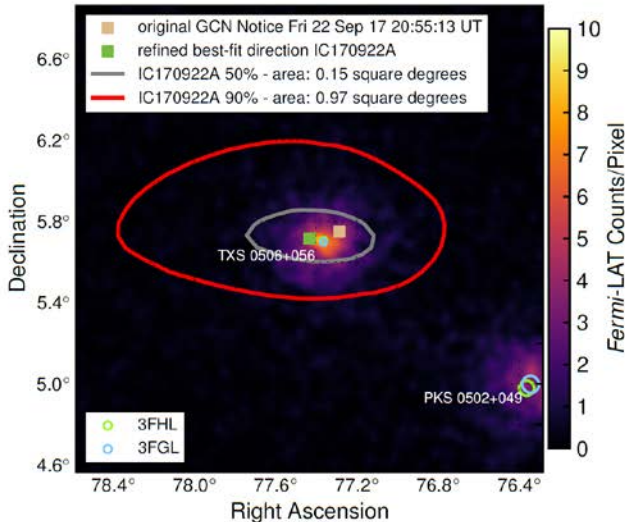
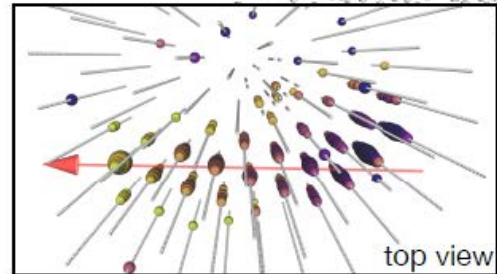
Illustrations: Katz et al. (2012)

Neutrino Production in Blazars

side view



IceCube-170922A



TXS 0506+056

(IceCube et al. 2018)

Basics of Neutrino Production

- $p + \gamma \rightarrow p + \pi^0$

or $n + \pi^+$

$$\pi^0 \rightarrow 2\gamma$$

$$\pi^+ \rightarrow \mu^+ + \nu_{\mu} \quad \tau = 2.55 \times 10^{-8} \text{ s}$$

$$\pi^- \rightarrow \mu^- + \nu_{\mu}$$

$$\mu^+ \rightarrow e^+ + \nu_{\mu} + \nu_e \quad \tau = 2.2 \times 10^{-6} \text{ s}$$

$$\mu^- \rightarrow e^- + \nu_{\mu} + \nu_e$$

$p\gamma$ likely strongly dominant over pp
in AGN jet environments.

Photo-Pion Production

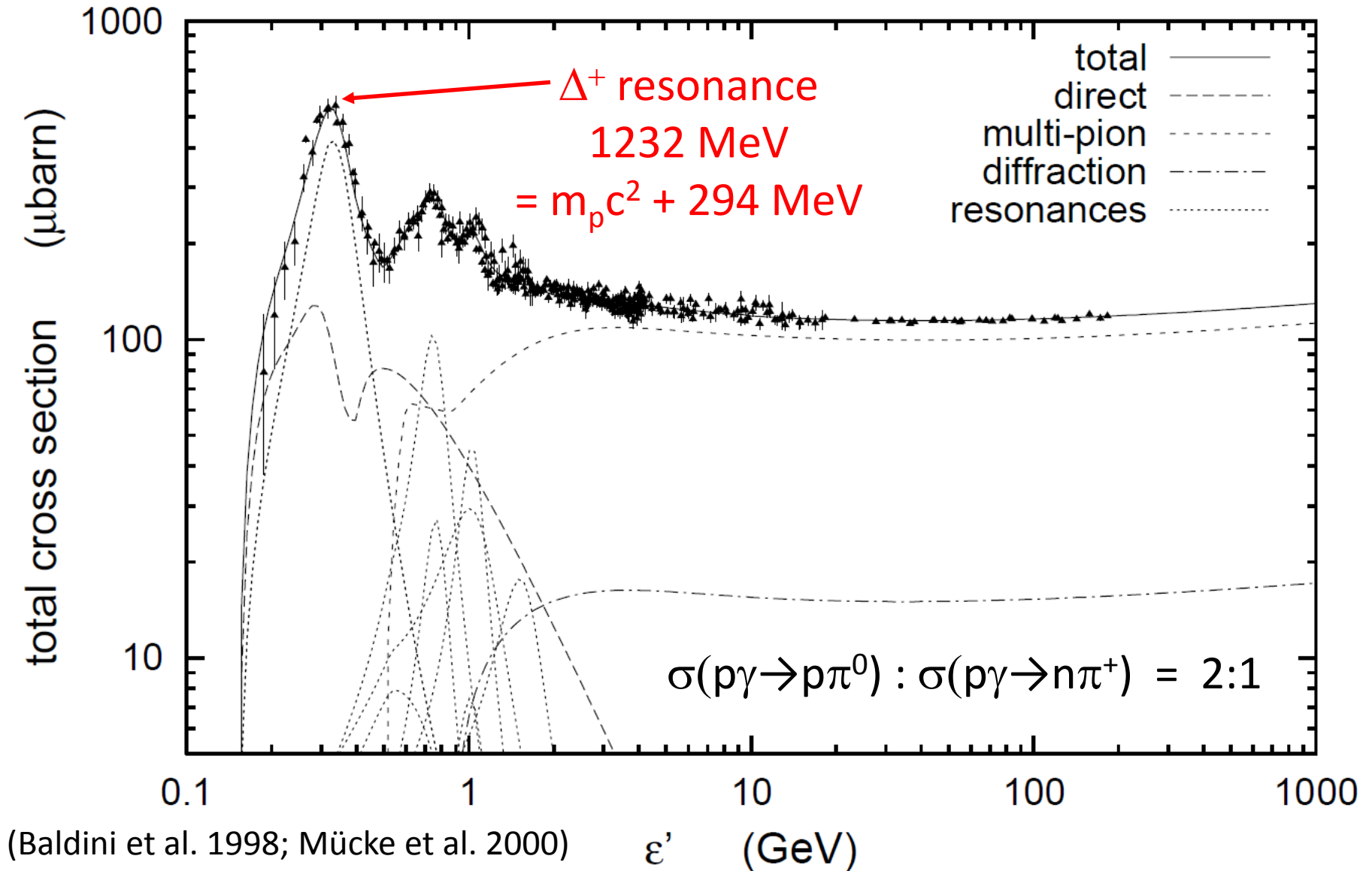
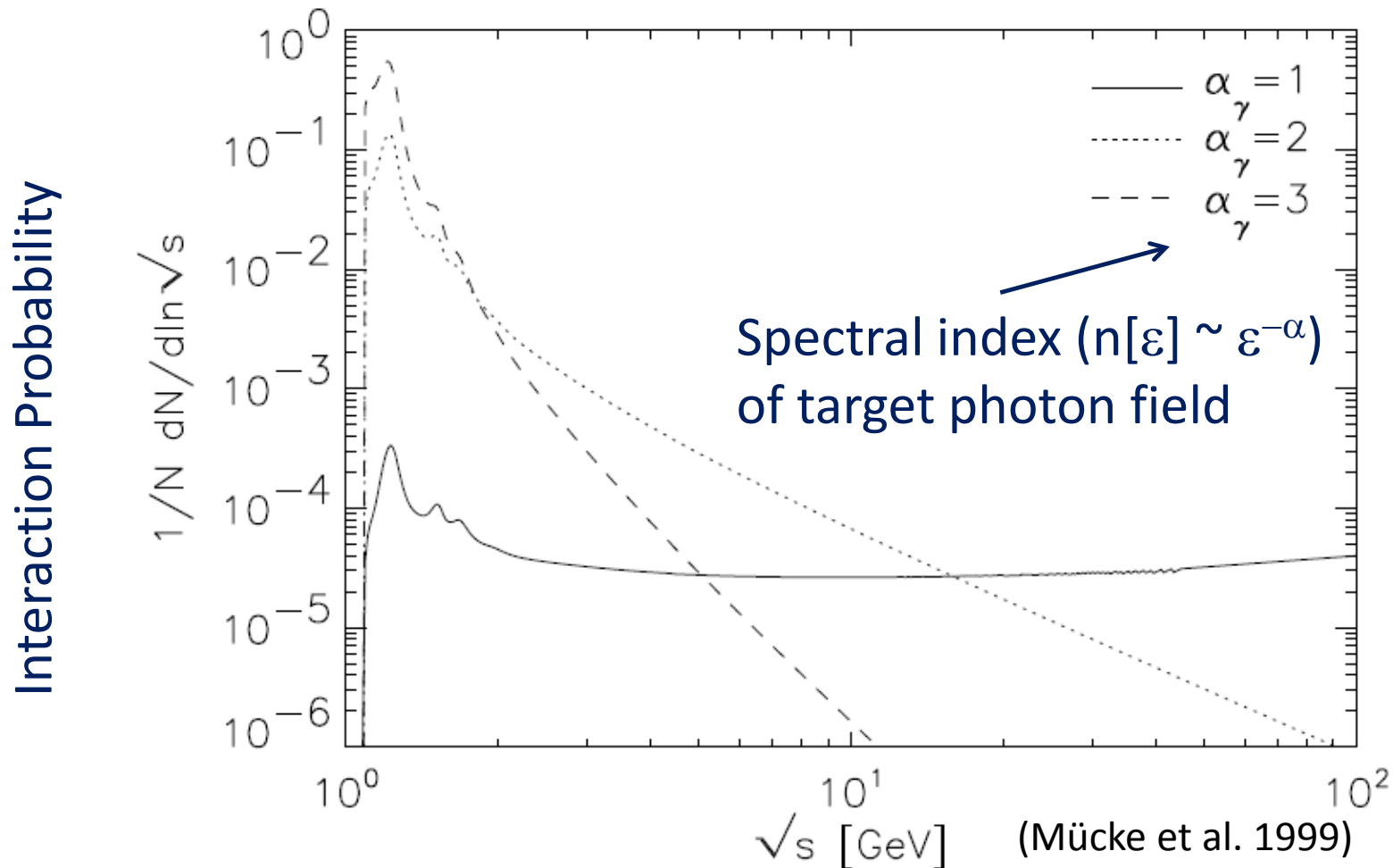


Photo-Pion Production



Center-of-Momentum energy

For realistic target photon fields,
most interactions occur near Δ^+ resonance.

Photo-pion production - Energetics

At Δ^+ resonance:

$$s = E'_p E'_t (1 - \beta_p' \mu) \sim E'_p E'_t \sim E_{\Delta^+}^2 = (1232 \text{ MeV})^2$$

and

$$E'_\nu \sim 0.05 E'_p$$

\Rightarrow To produce IceCube neutrinos ($\sim 100 \text{ TeV} \rightarrow E_\nu = 10^{14} E_{14} \text{ eV}$):

$$\text{(i.e., } E'_\nu = 10 E_{14} \delta_1^{-1} \text{ TeV)}$$

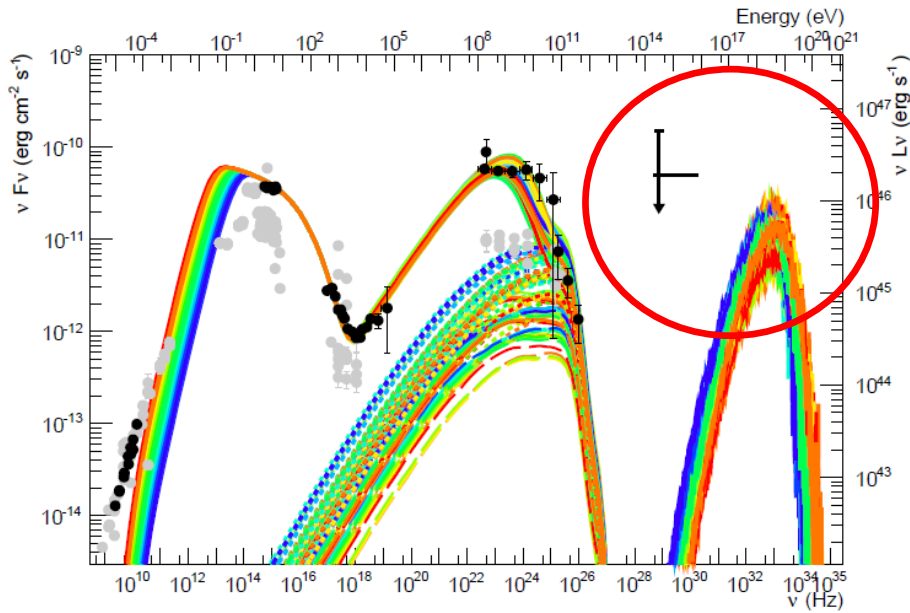
Need protons with

$$E'_p \sim 200 E_{14} \delta_1^{-1} \text{ TeV} \Rightarrow \text{PeV CRs}$$

and target photons with

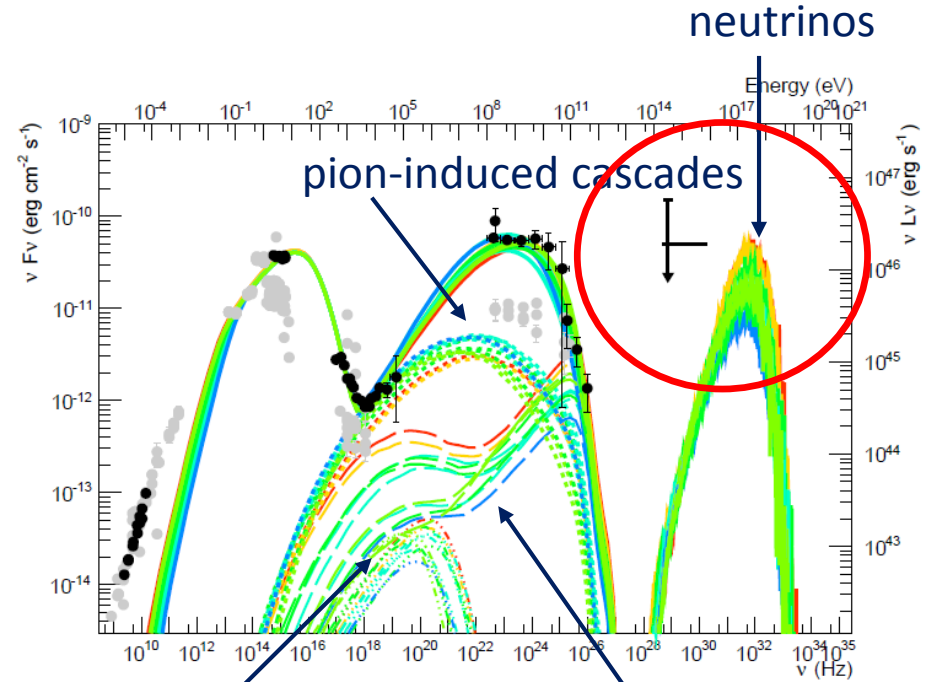
$$E'_t \sim 1.6 E_{14}^{-1} \delta_1 \text{ keV} \Rightarrow \text{X-rays}$$

Photo-Pion Models for TXS 0506+056



(a) Proton synchrotron modeling of TXS 0506+056

(Cerruti et al. 2019: MNRAS, 483, L12)



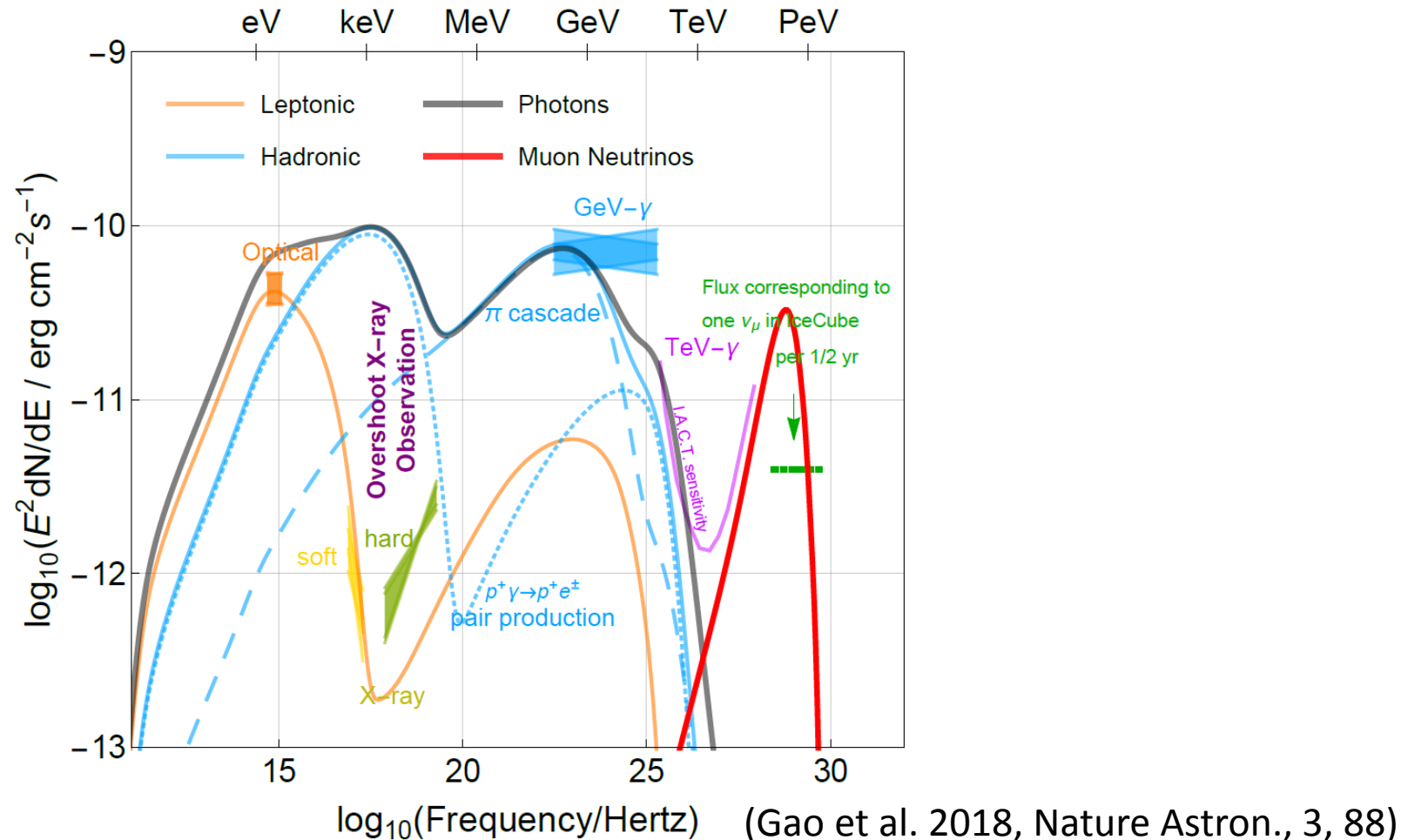
(b) Lepto-hadronic modeling of TXS 0506+056

Proton-synchrotron

Bethe-Heitler-induced
cascades

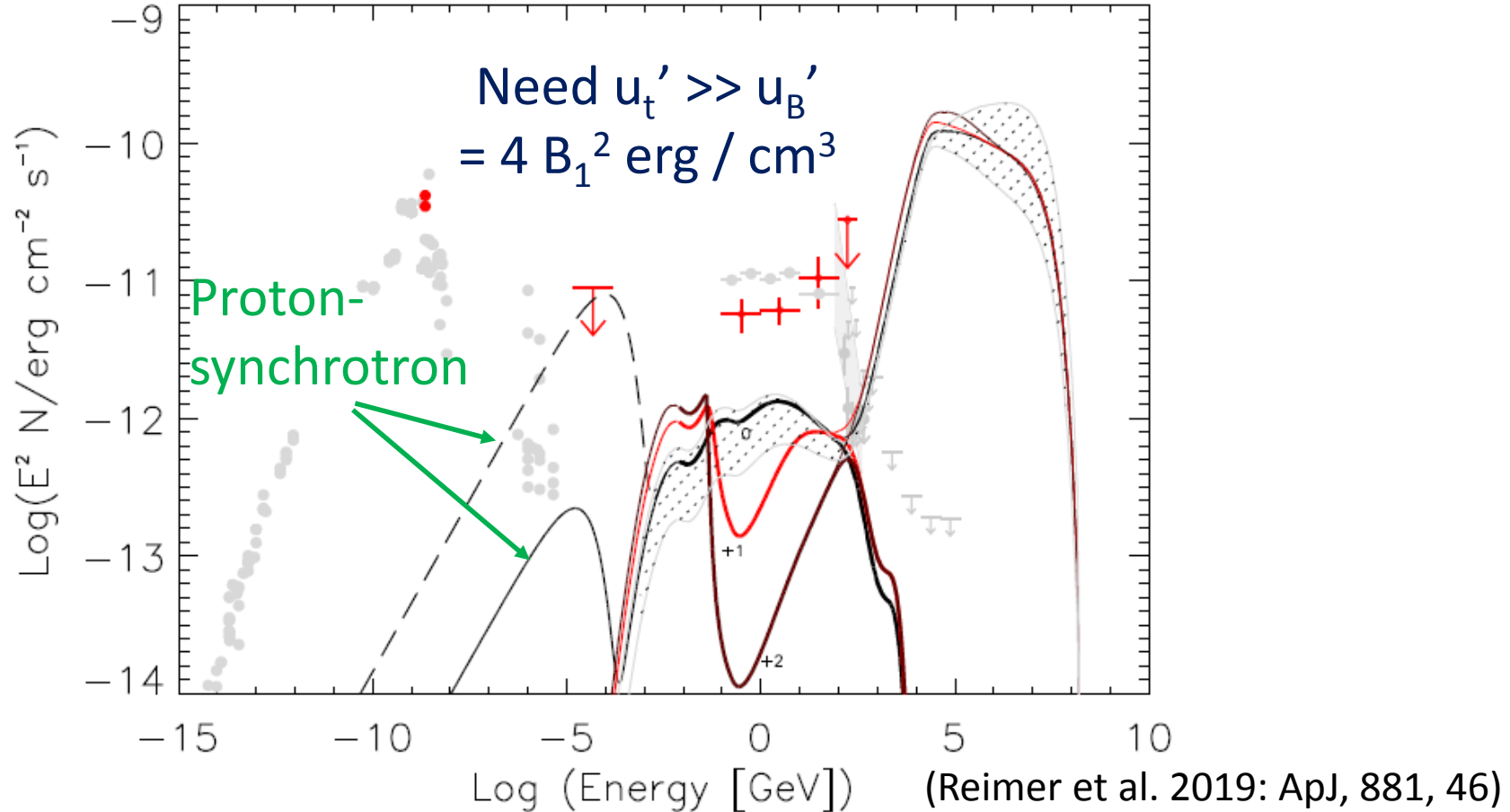
Models producing neutrinos and gamma-rays
through the same proton population, predict too
high neutrino energies!

Photo-Pion Models for TXS 0506+056



Models with $p\text{-}\gamma$ induced γ -ray emission over-produce X-rays due to cascades!

Photo-Pion Models for TXS 0506+056

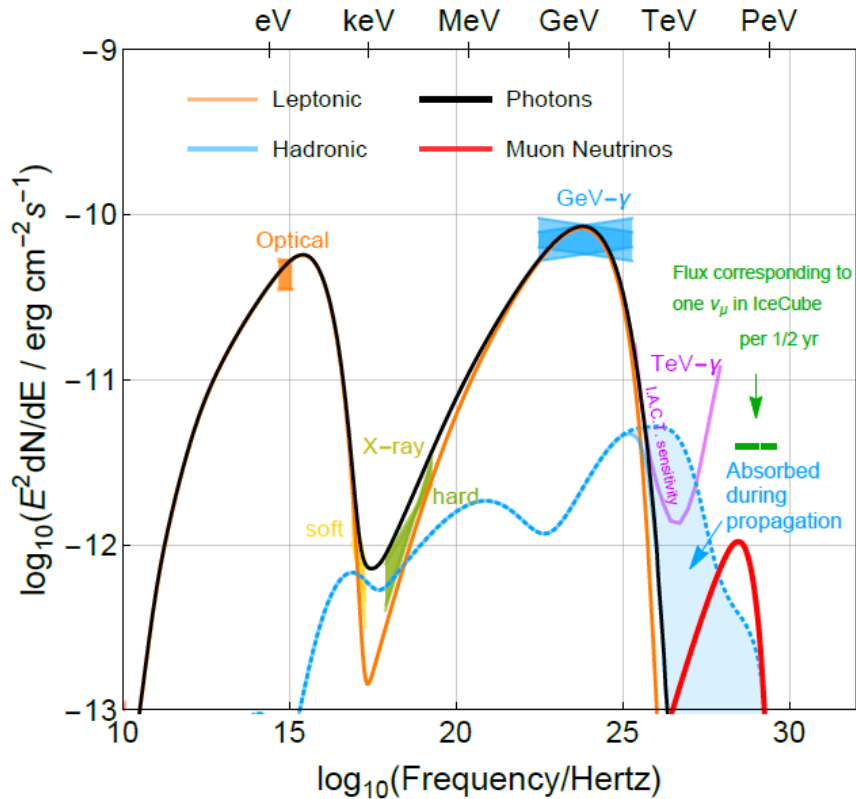


From cascading constraints:

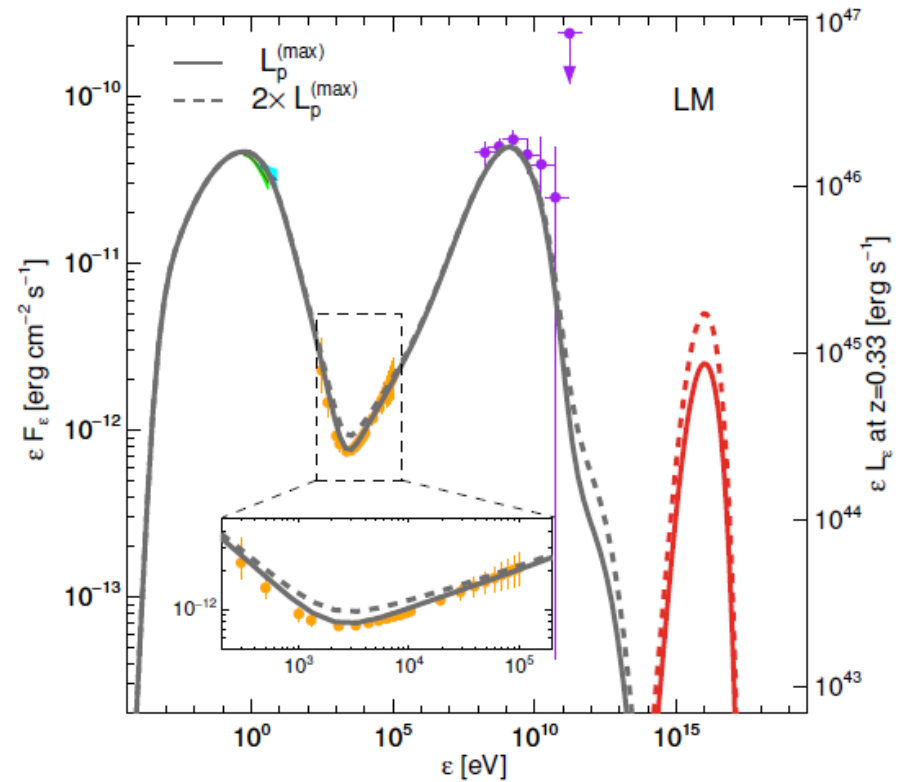
$\text{p}\gamma$ neutrino production in TXS 0506+056 possible with strong external UV/X-ray radiation field, but under-predicts Fermi γ -rays.

=> No neutrino – γ -ray correlation expected!

Photo-Pion Models for TXS 0506+056



(Gao et al. 2018, Nature Astron., 3, 88)

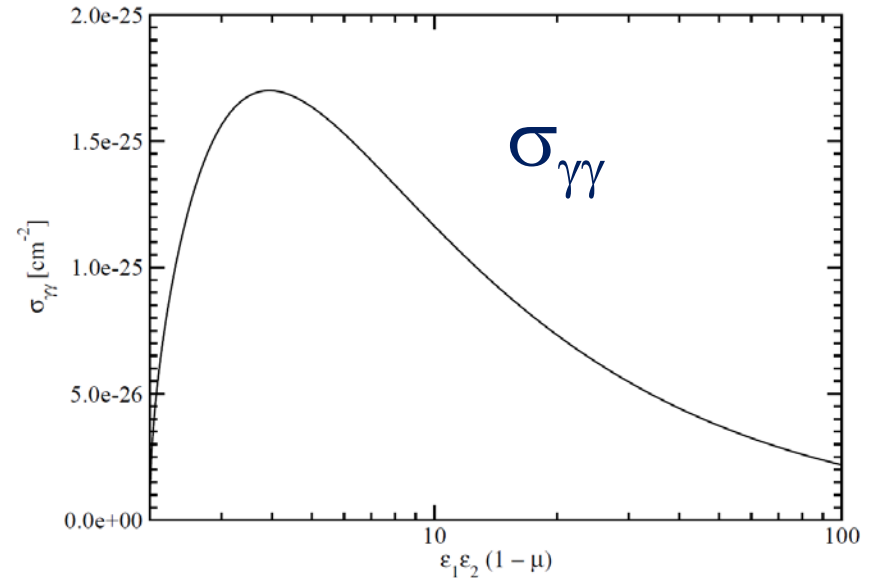
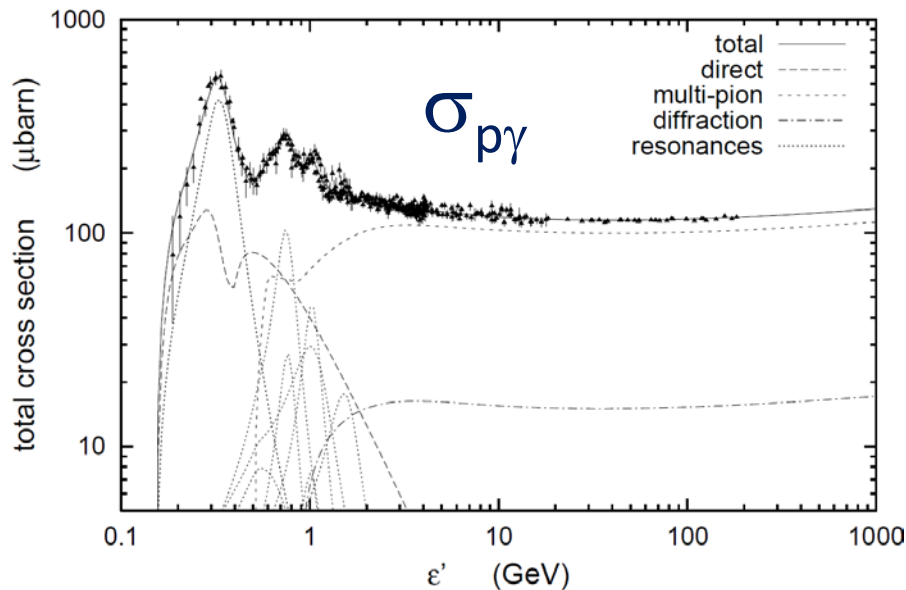


(Keivani et al. 2018, ApJ, 864, 84)

Models producing neutrinos and gamma-rays require leptonically dominated gamma-ray production!

The $p\gamma$ Efficiency Problem

- Efficiency for protons to undergo $p\gamma$ interaction $\sim \tau_{p\gamma} = \ell_{\text{esc}} \sigma_{p\gamma} n_{\text{ph}}$
- Likelihood of γ -ray photons to be absorbed $\sim \tau_{\gamma\gamma} = R \sigma_{\gamma\gamma} n_{\text{ph}}$



$$\frac{\tau_{p\gamma}}{\tau_{\gamma\gamma}} = \frac{\sigma_{p\gamma} \ell_{\text{esc}}}{\sigma_{\gamma\gamma} R} \approx \frac{1}{300} \frac{\ell_{\text{esc}}}{R}$$

ℓ_{esc} = average length
travelled by protons
until escape

at $E_\gamma \sim \frac{m_e^2 c^4}{E_t} \sim 3.3 \times 10^{-5} E_\nu \longleftarrow \sim \text{GeV} - \text{TeV } \gamma\text{-rays}$

The $p\gamma$ Efficiency Problem

$$\frac{\tau_{p\gamma}}{\tau_{\gamma\gamma}} = \frac{\sigma_{p\gamma} \ell_{\text{esc}}}{\sigma_{\gamma\gamma} R} \approx \frac{1}{300} \frac{\ell_{\text{esc}}}{R}$$

ℓ_{esc} = average length travelled by protons until escape

ℓ_{esc} from random walk:

mean free path $\lambda = \eta(\gamma) r_g(\gamma)$

Number of scatterings to escape, N_s : $R = \sqrt{N_s} \lambda$

$$\ell_{\text{esc}} = N_s \lambda = \frac{R^2}{\lambda} \approx 3.3 \times 10^{21} \eta(\gamma)^{-1} R_{16}^2 B_2 E_{15}^{-1} \text{ cm}$$

$$\Rightarrow \frac{\tau_{p\gamma}}{\tau_{\gamma\gamma}} = \frac{\sigma_{p\gamma} \ell_{\text{esc}}}{\sigma_{\gamma\gamma} R} \approx 1.1 \times 10^3 \eta(\gamma)^{-1} R_{16} B_2 E_{15}^{-1}$$

\Rightarrow Proton $p\gamma$ efficiency can be $\gg \tau_{\gamma\gamma}$, but misleading, as $t_{\text{cool},p\gamma}$ and $t_{\text{esc},p} \gg R/c$

The $p\gamma$ Efficiency Problem

Relevant constraint for proton bulk kinetic luminosity:

$$L'_\nu \approx \frac{1}{2} mpc^2 \int d\gamma_p Np(\gamma_p) |\dot{\gamma}_{p,p\gamma}| = L'_\nu (\text{obs})$$

$$\dot{\gamma}_{p,p\gamma} \approx -c \langle \sigma_{p\gamma} f \rangle \frac{u'_{ph}}{m_e c^2} \gamma_p$$

$$\Rightarrow L_p u'_{ph} \approx 1.4 \times 10^{52} \delta_1^{-4} \Gamma_1^2 R_{16}^{-1} \left(\frac{\text{erg}}{\text{s}} \right) \left(\frac{\text{erg}}{\text{cm}^3} \right)$$

(Reimer et al. 2019: ApJ, 881, 46)

Cosmic-Ray Acceleration in Blazars

- No conclusive correlation between AGN and UHECR arrival directions observed.
- To produce > 100 TeV neutrinos \rightarrow Need PeV protons
- Simplest constraint: Confinement (Hillas Criterion):

$$E_n < Z e B R = 3 \times 10^{18} B_2 R_{16} \text{ eV}$$

($Z = 1$ for protons)

for hadronic blazar models with $B = 100 B_2 \text{ G}$

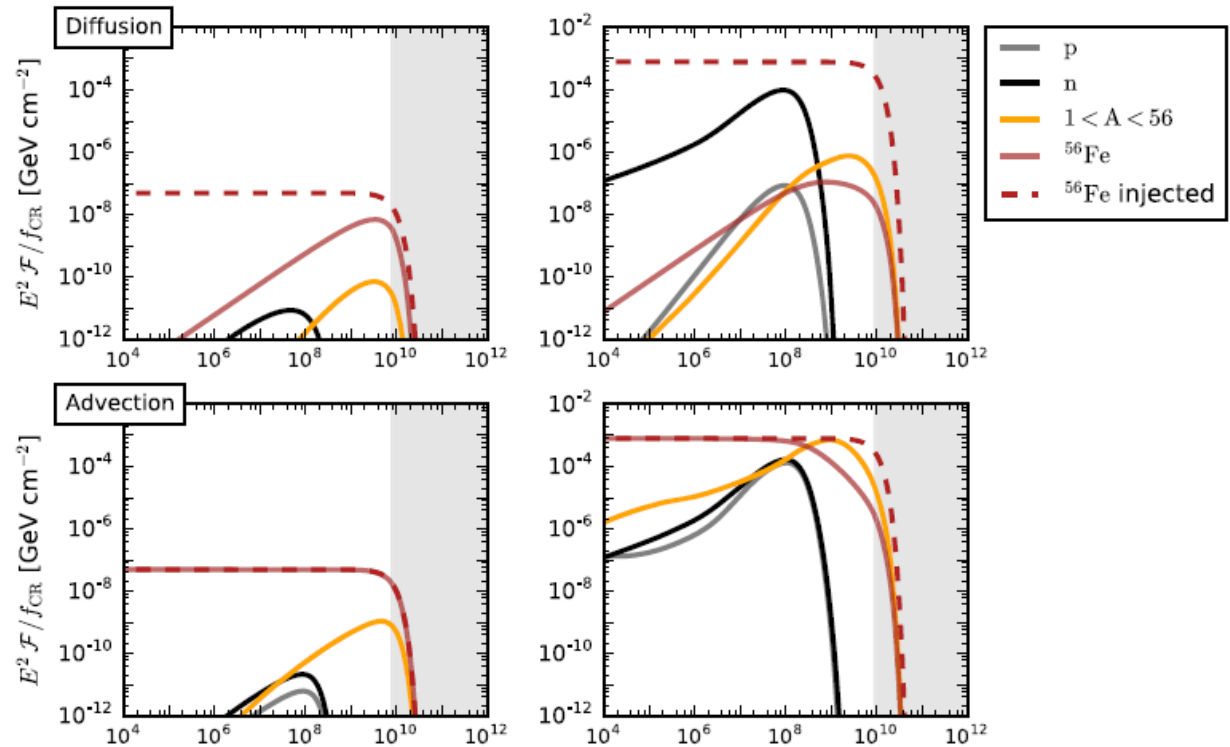
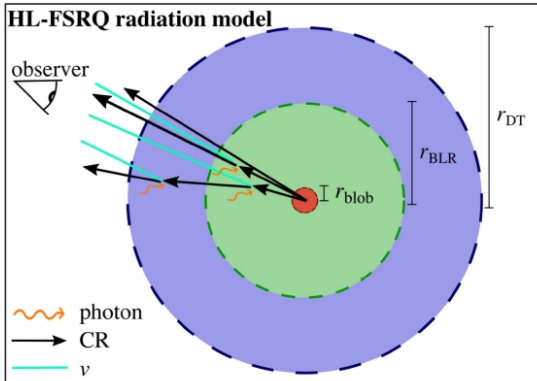
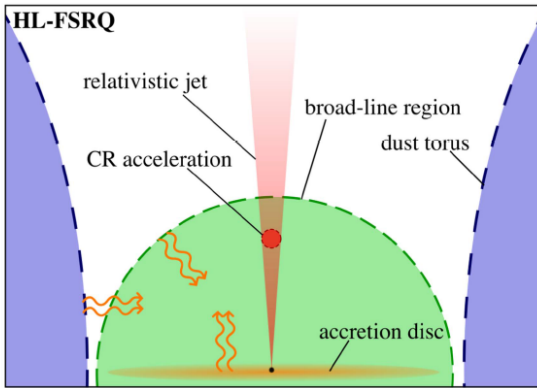
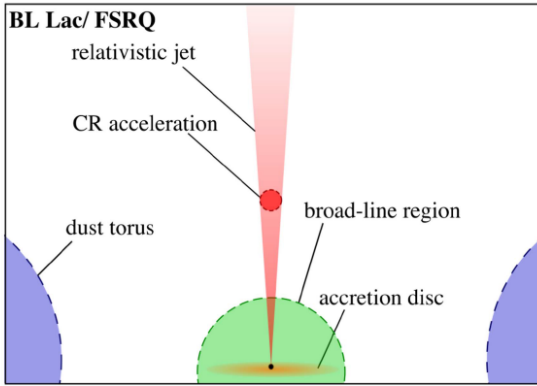
$$R = 10^{16} R_{16} \text{ cm}$$

\Rightarrow Protons for IceCube neutrino production: 

\Rightarrow UHECRs ($E > 10^{19}$ eV): Plausible for heavy elements
($Z \gg 1$)

(e.g., Rodrigues et al., 2018: ApJ, 854, 54)

Cosmic-Ray Acceleration in Blazars



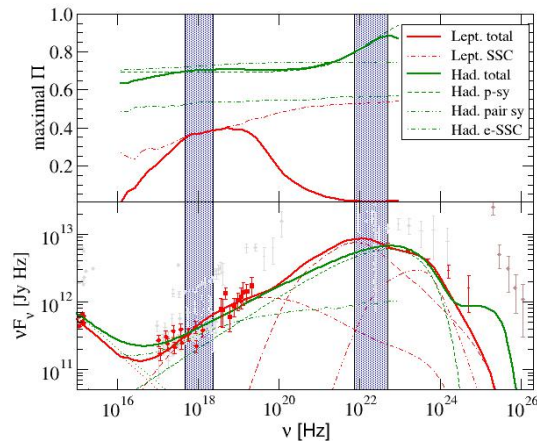
Acceleration of Fe and subsequent spallation may produce UHECR spectra and composition similar to the observed ones.

(Rodrigues et al., 2018: ApJ, 854, 54)

Observable signatures of proton / CR acceleration in blazars

- X-ray / γ -ray polarization

3C279

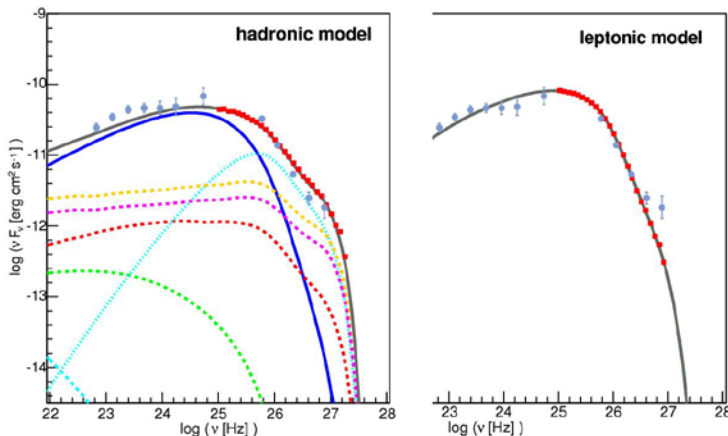


→ IXPE / AMEGO

Requires simultaneous optical polarimetry for comparison.

(Zhang & Böttcher 2013: ApJ, 744, 18)

- Spectral signatures of $p\gamma$ induced cascades



→ CTA

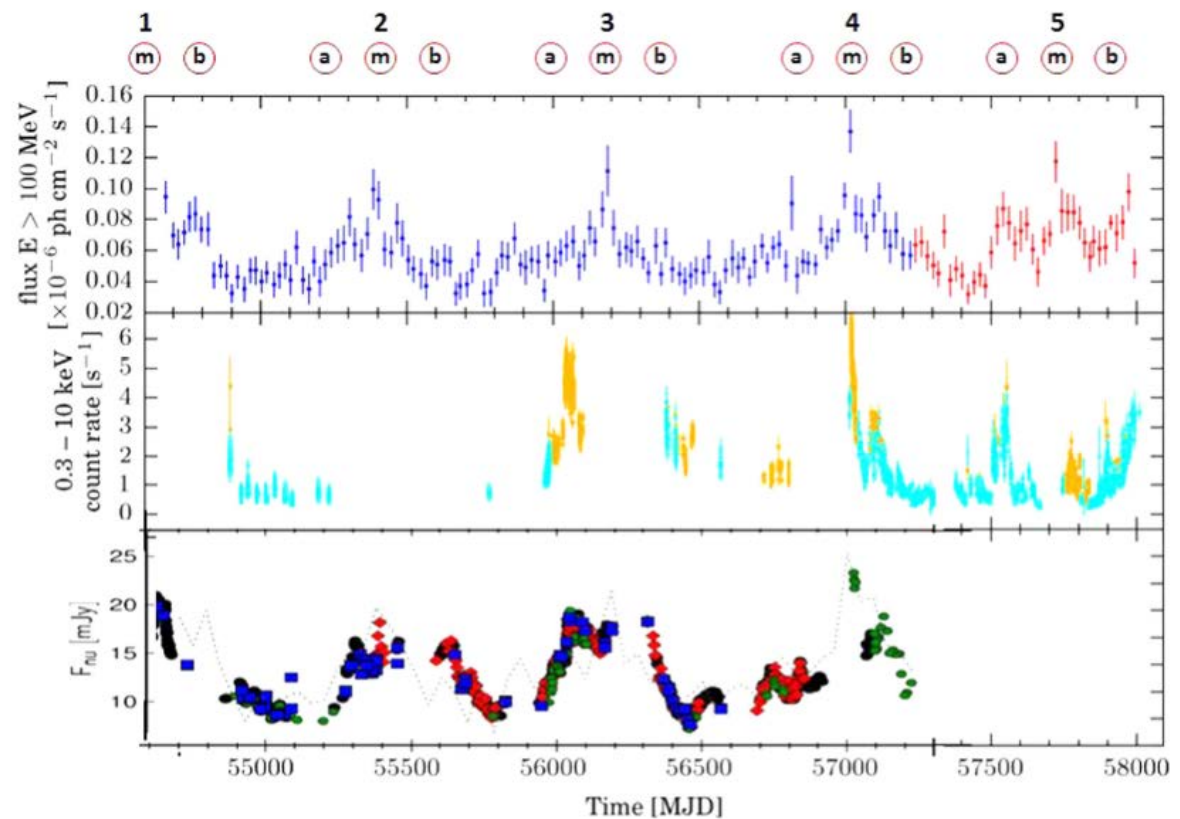
Requires simultaneous MWL (OIR-UV-X-ray) coverage.

(Zech et al. 2017: A&A, 602, A25)

Binary Super-Massive Black Holes – Gravitational Waves

Evidence for periodicity in some AGN, e.g., OJ 287, PG 1553+113

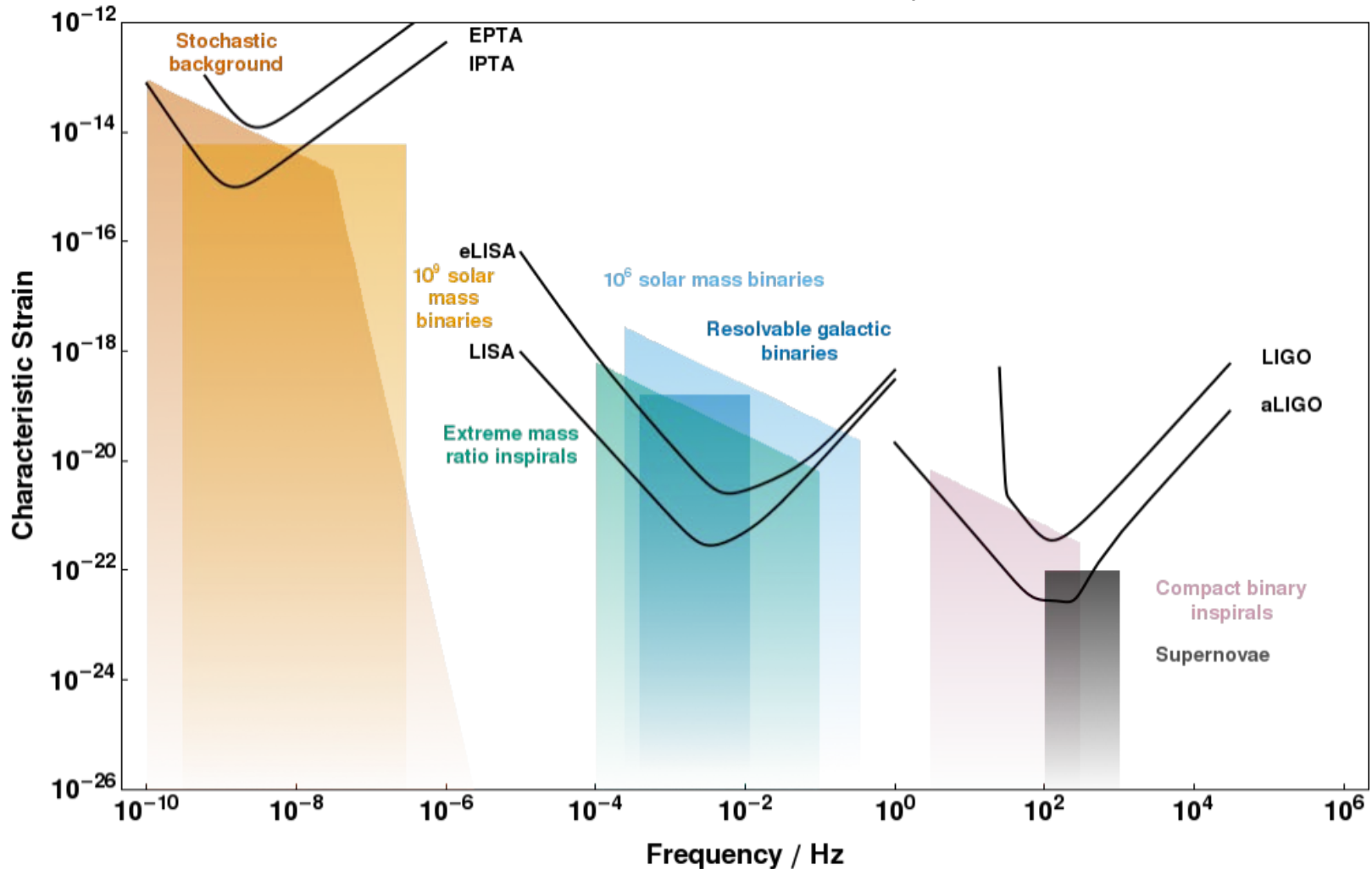
Possible signature
of binary super-
Massive black holes



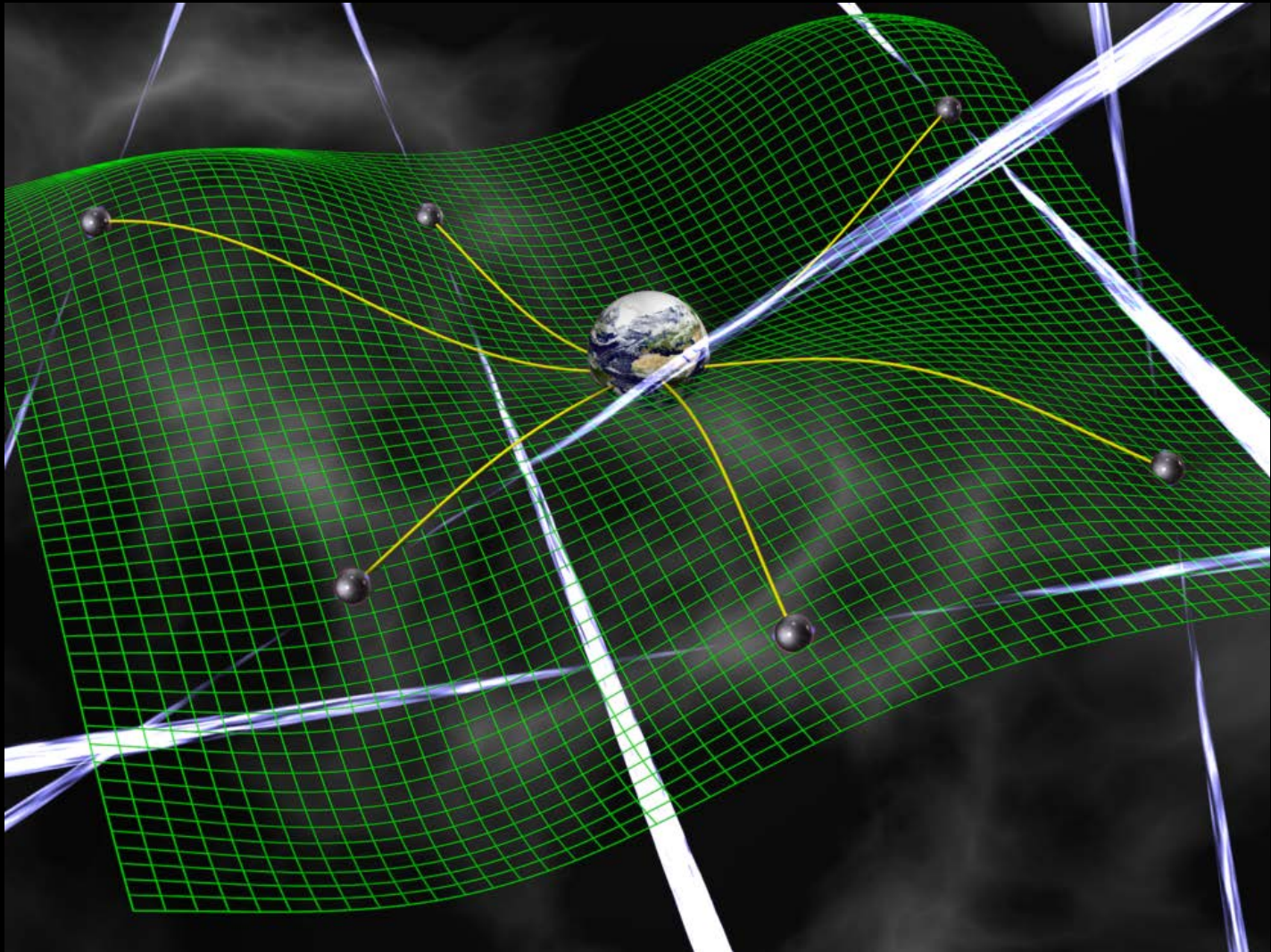
(Tavani et al., 2018, ApJ, 851, 11)

Gravitational Waves from supermassive black hole binaries

GW frequencies: $f \sim 30 P_y^{-1}$ nHz



Pulsar Timing Arrays

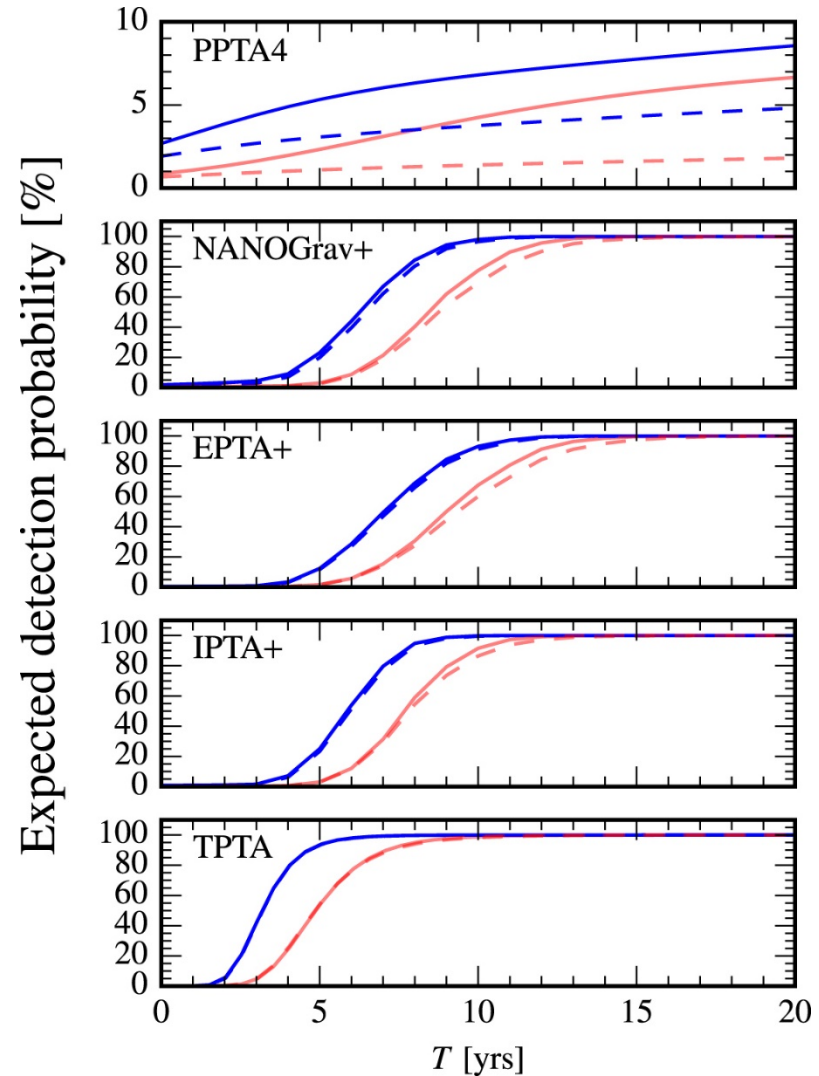


Search for relative distortions of pulsar timing signals
in various directions

Gravitational Waves from supermassive black hole binaries

Prospects for detection with
Pulsar Timing Arrays (PTAs):

Large Pulsar Timing Arrays
(incl. Parkes PTA) may detect
SMBH binary GW background
within the next $\sim 5 - 10$ years.



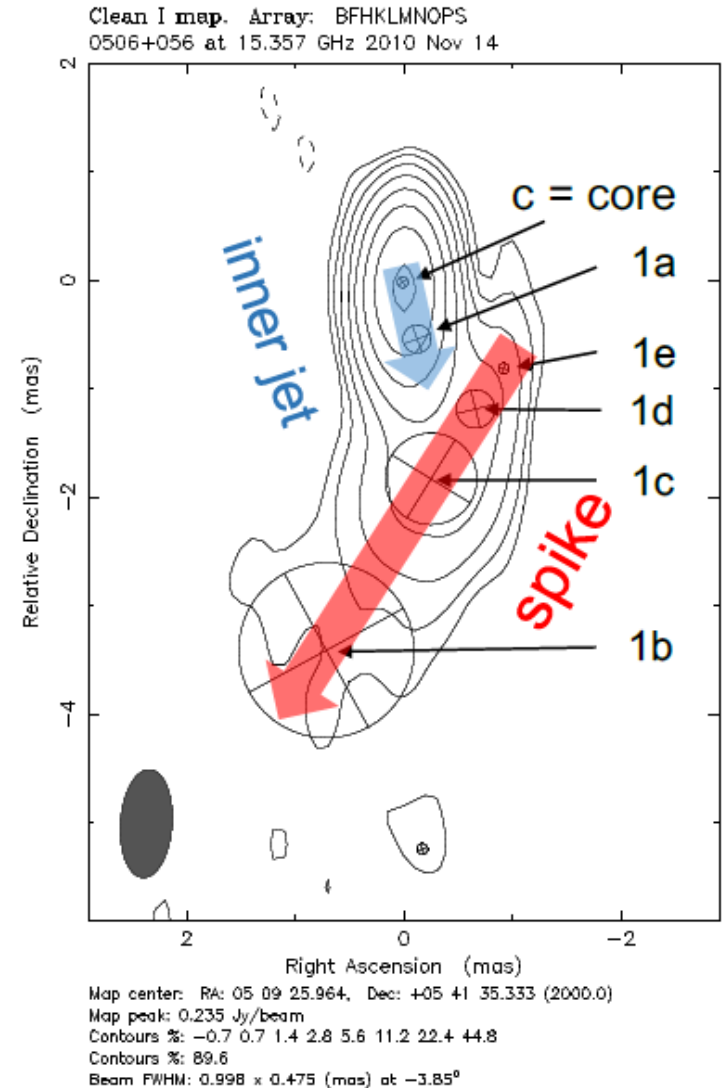
(Taylor et al. 2016, ApJ, 819, L6)

Supermassive black hole binaries

No detection of individual binaries, but estimates of total number of binary SMBHs.

- disk/jet precession
- feasibility of precessing-jet models for blazar variability

SMBH binaries: Possible site of jet-jet interaction → VHE neutrino production (?)



(Britzen et al., 2019: A&A, 630, 103)

Summary

- Production of IceCube neutrinos in AGN / blazars plausible, but no correlation between γ -ray and neutrino activity necessarily expected.
- Blazar jets are plausible accelerators of PeV – EeV CRs; UHECR acceleration plausible for primary acceleration of heavier nuclei (Fe).
- Binary Supermassive black holes (Periodicity / precessing / interacting jets) produce nHz gravitational-wave background, which may be detectable by PTAs within the next decade.



Thank you!

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