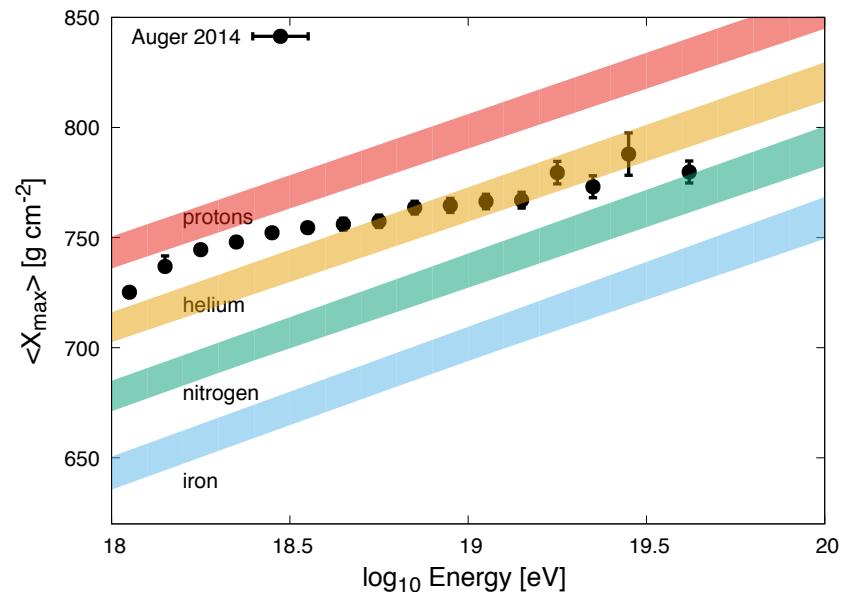
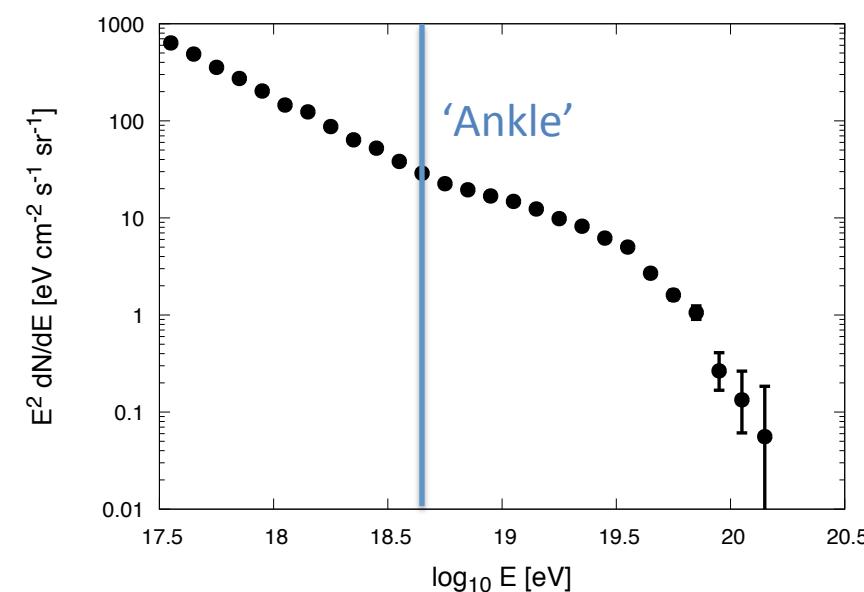
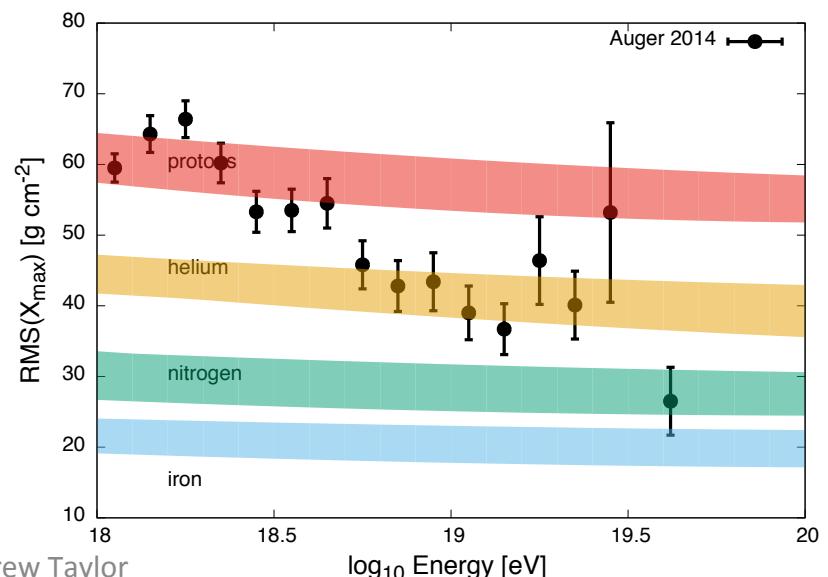
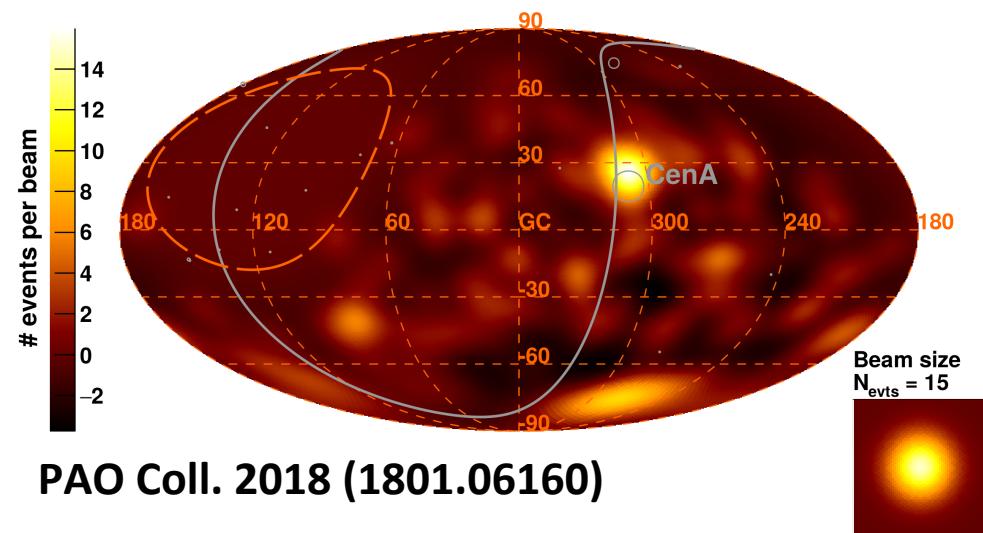


CTA as a Probe of Extragalactic CR Sources



Observed Excess Map - $E > 60$ EeV



CTA Symposium- Andrew Taylor

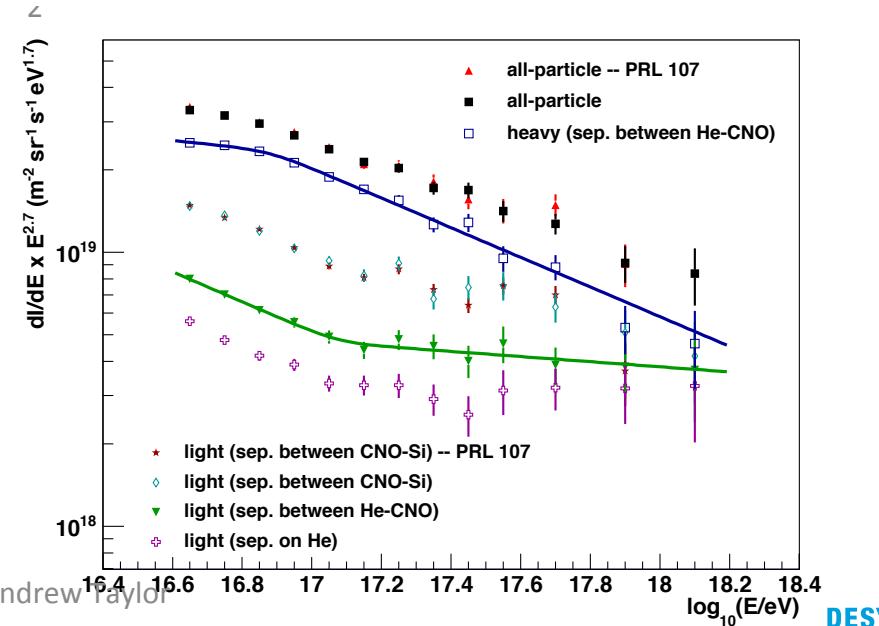
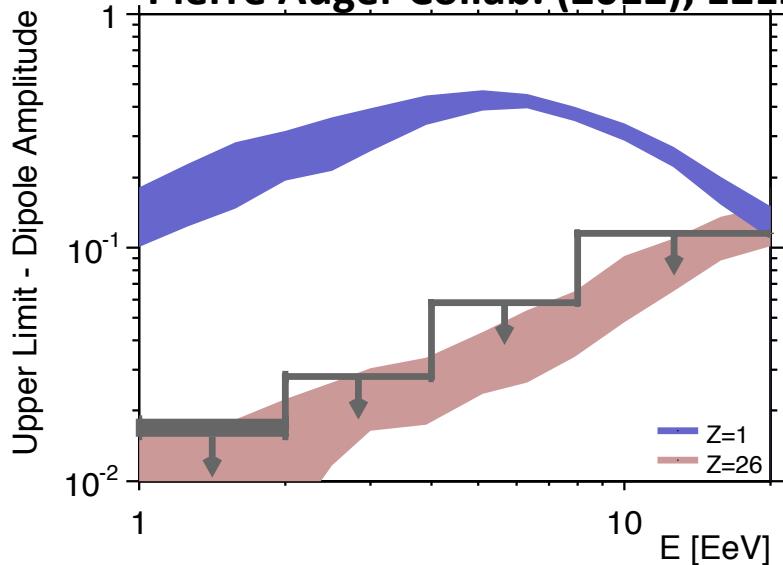
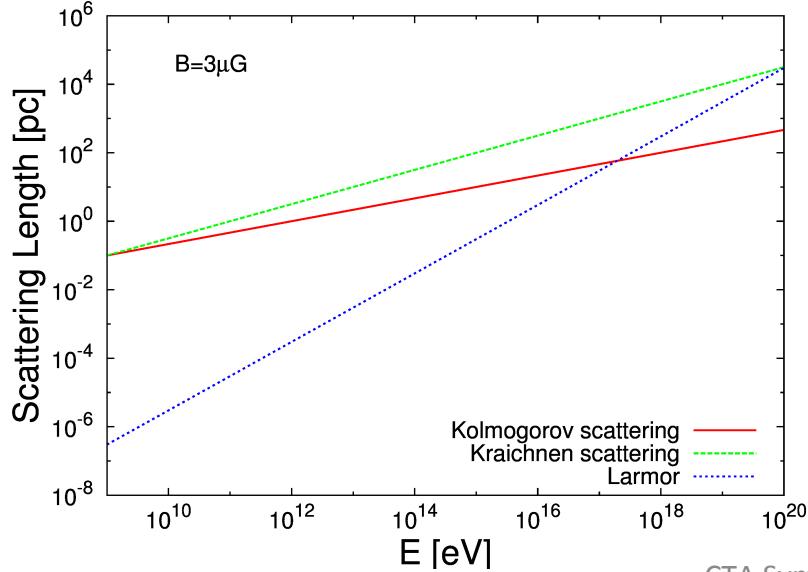
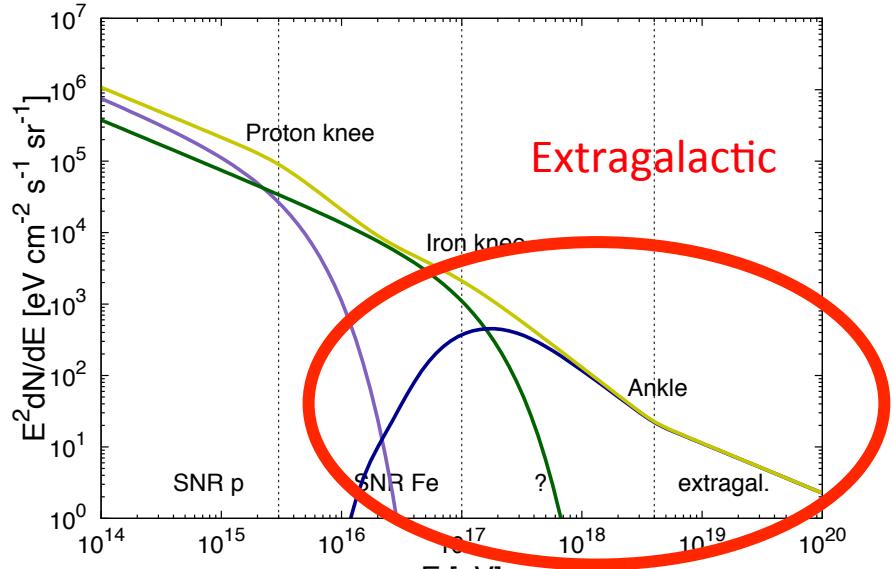
By Andrew Taylor

Where Does Extragalactic Begin?

Anisotropy constraint:

Giacinti et al. (2011), 1112.5599

Pierre Auger Collab. (2012), 1212.3083



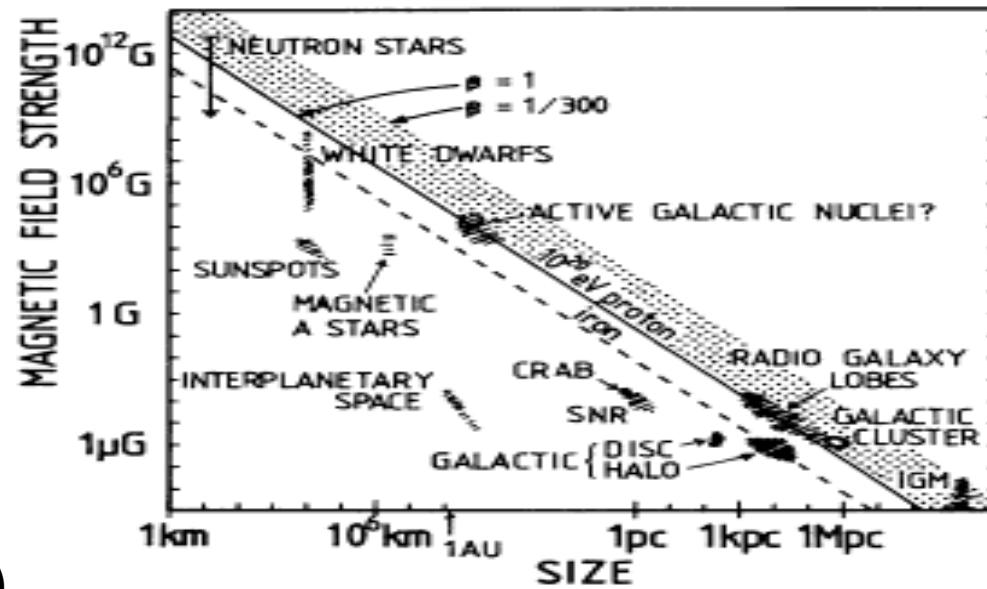
Candidate Sources- Basic Argument Hasn't Changed!

$$t_{\text{acc}} = \eta \frac{R_{\text{lar}}}{c\beta^2}$$

$$R_{\text{lar}} = \frac{\beta}{\eta} R$$

$$t_{\text{esc.}} = \frac{R^2}{\eta c R_{\text{lar}}}$$

$\eta \approx 1$ assumed in plot below



AM Hillas (1984)

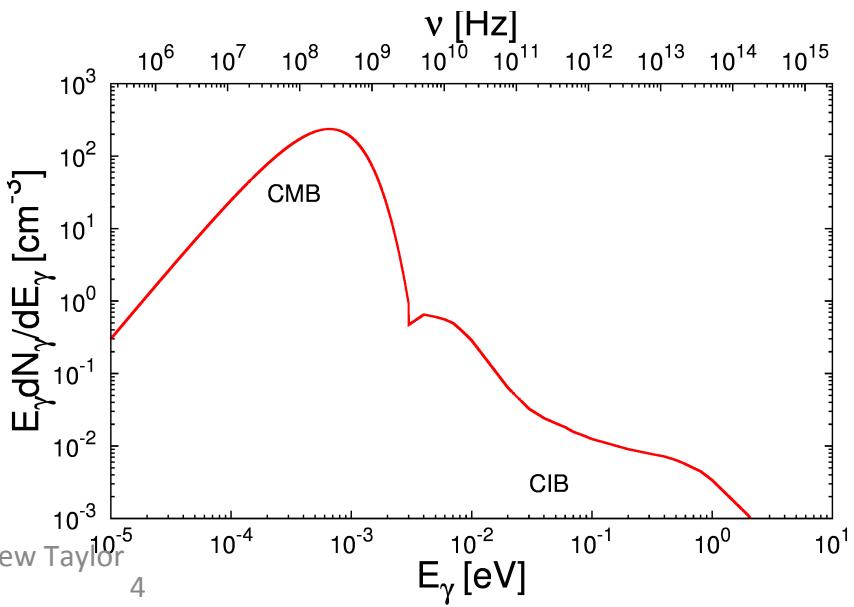
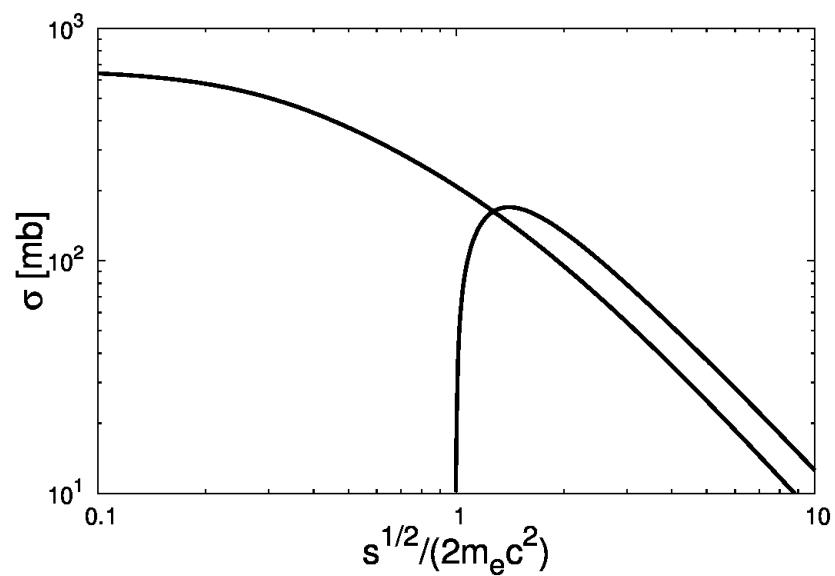
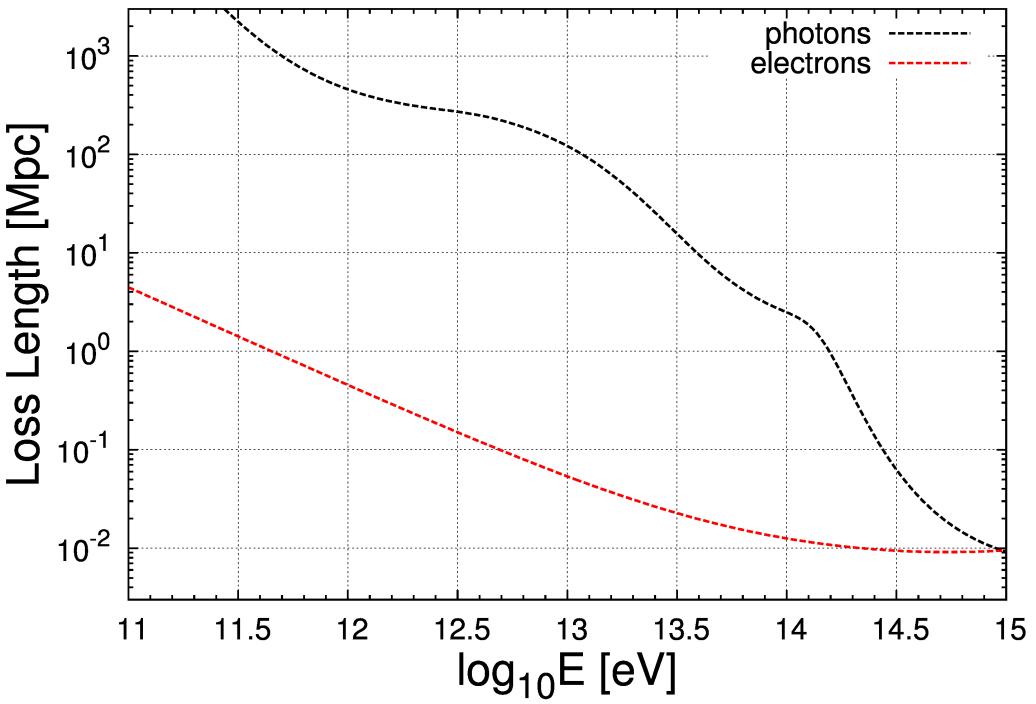
$$L > 3 \times 10^{42} \frac{1}{\beta^2} \left(\frac{E_p}{3 \times 10^{18} \text{ eV}} \right)^2 \text{ erg s}^{-1}$$

AGN

GRB

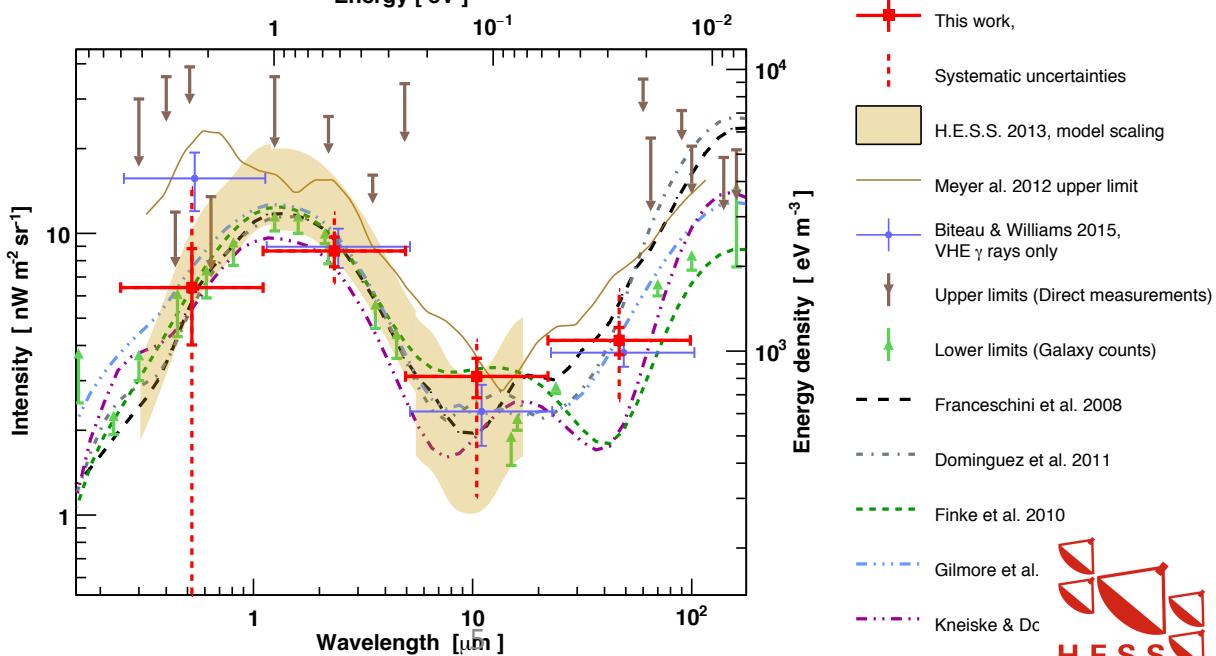
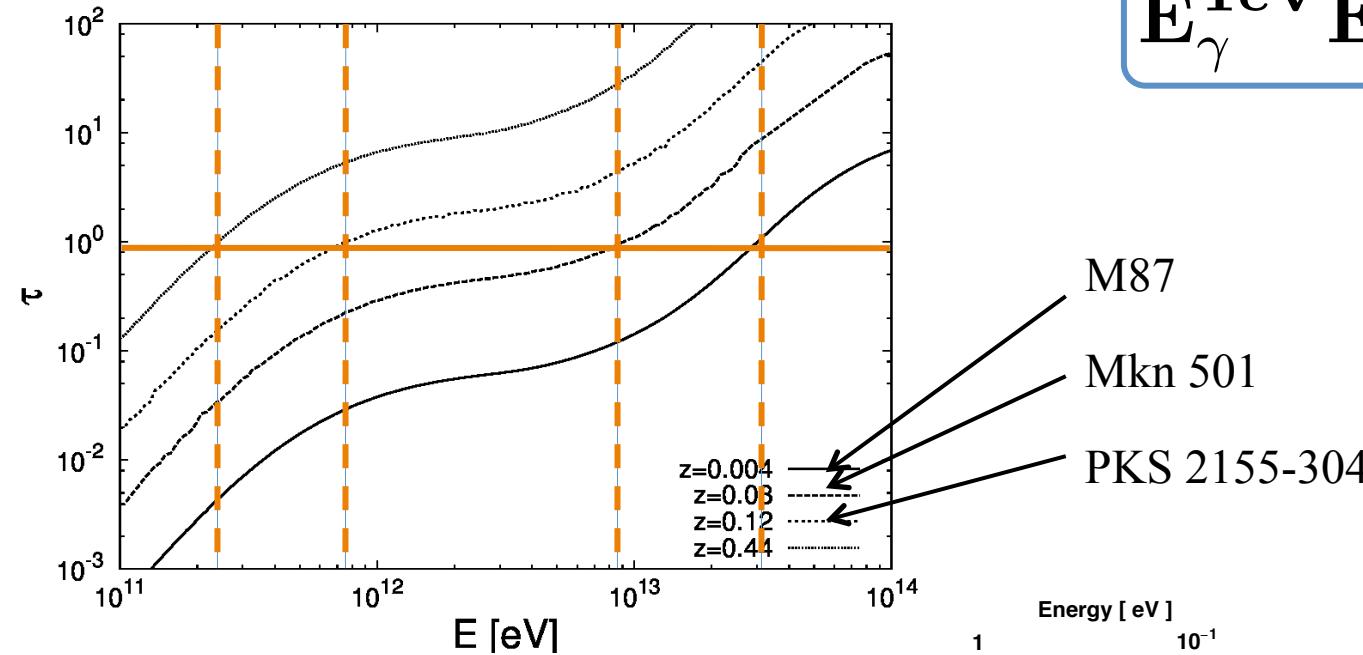
CTA Symposium- Andrew Taylor

Gamma-Ray Blazars (AGN) and the EBL



EBL Attenuation

$$E_\gamma^{\text{TeV}} E_\gamma^{\text{eV}} \approx 1$$

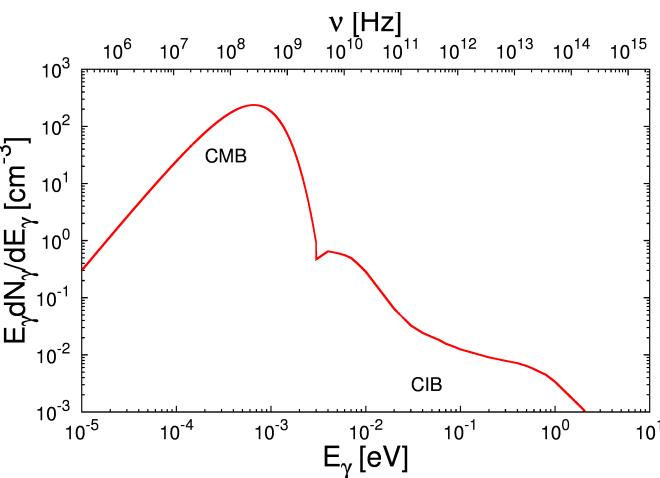
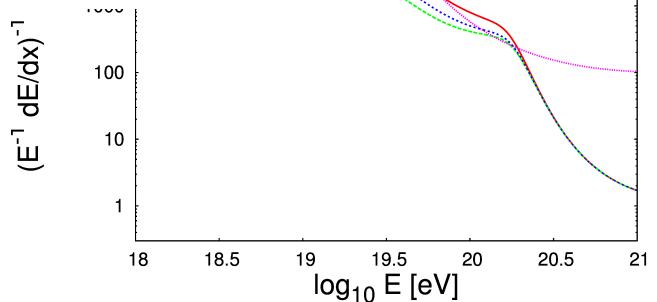
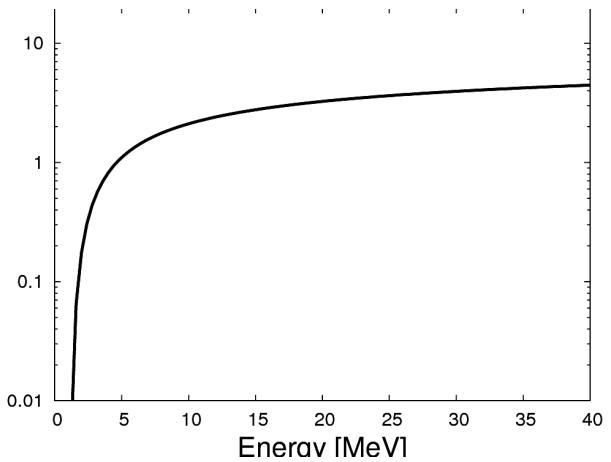
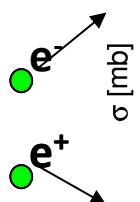
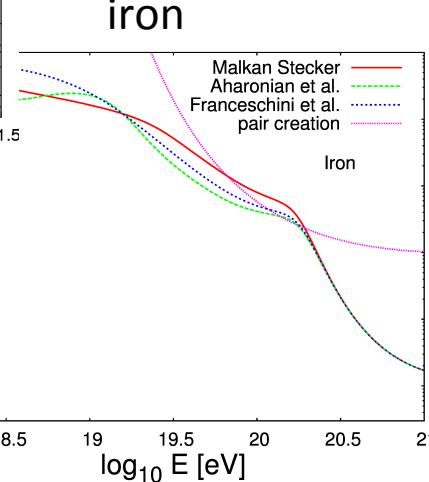
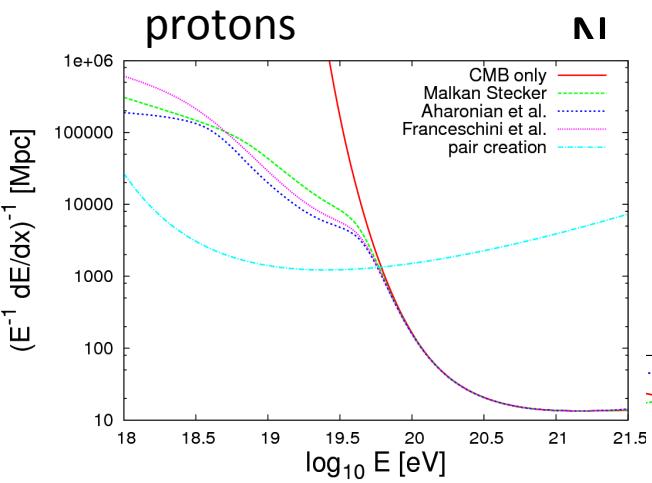
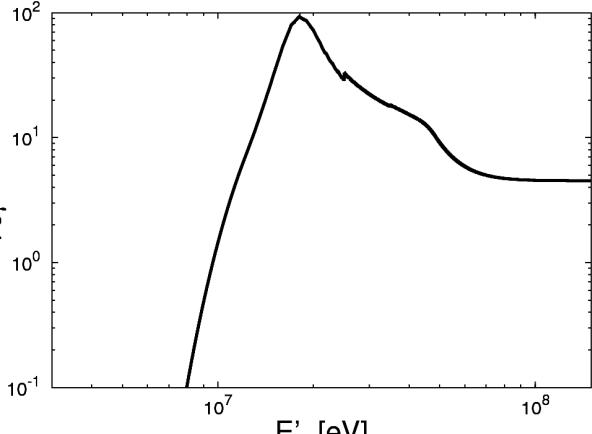


HESS Coll. A&A 606 (2017) A59

CTA Symposium- Andrew Taylor



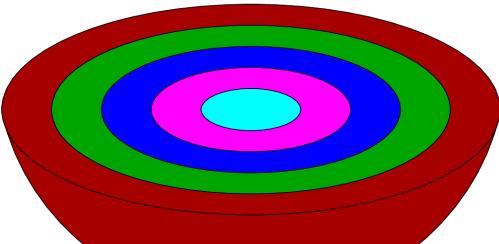
Cosmic Ray Proton/Nuclei Interactions



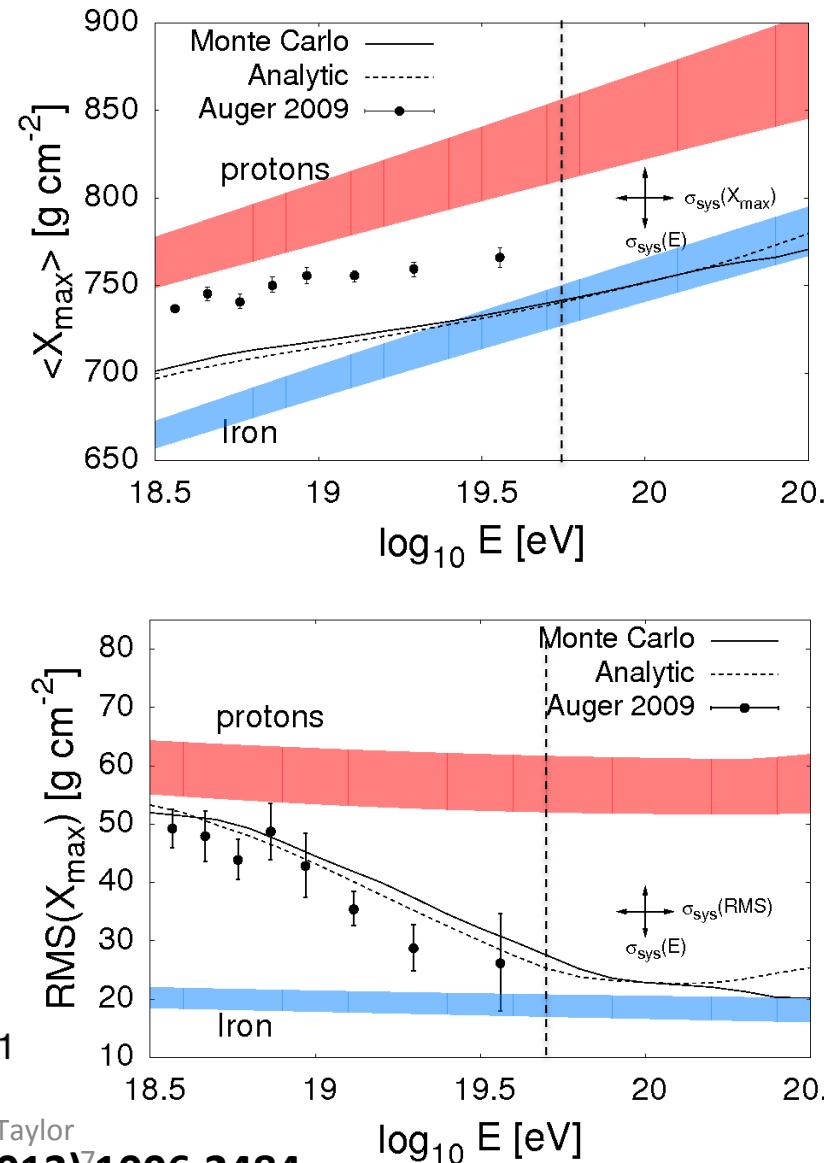
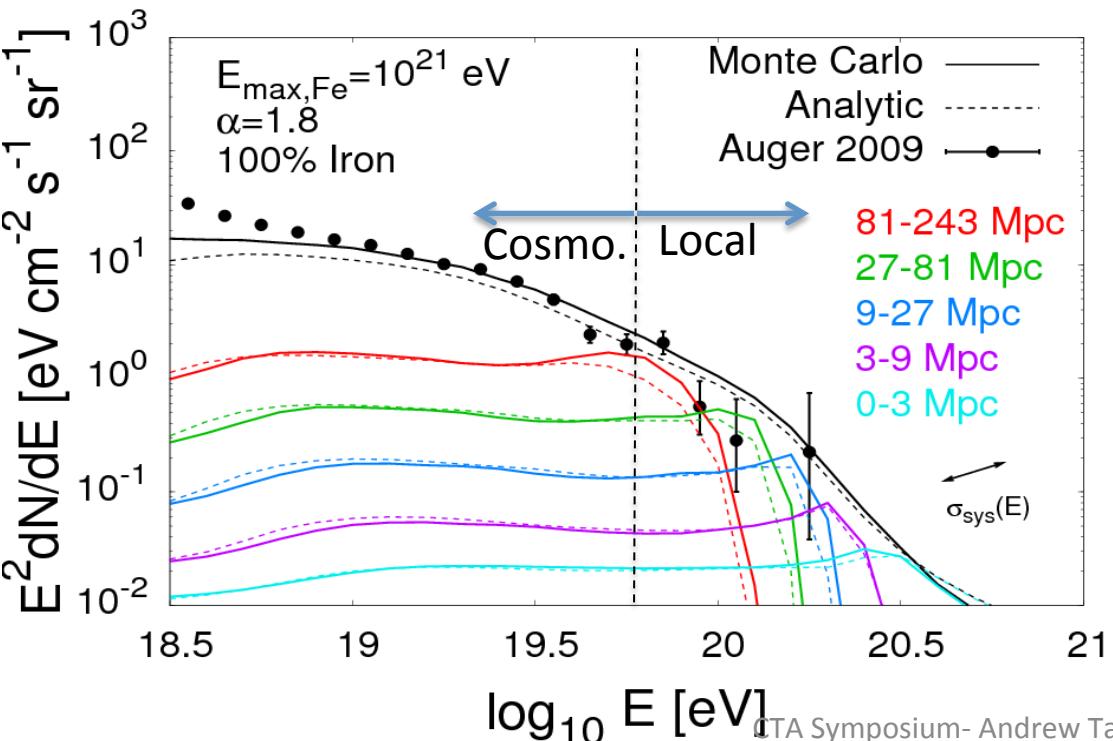
$$E_{\text{CR}}^{30\text{PeV}} E_{\gamma}^{\text{eV}} \approx A$$

Local Scales Effect Highest Energies: Analytic Treatments

0 3 9 27 81 243 Mpc



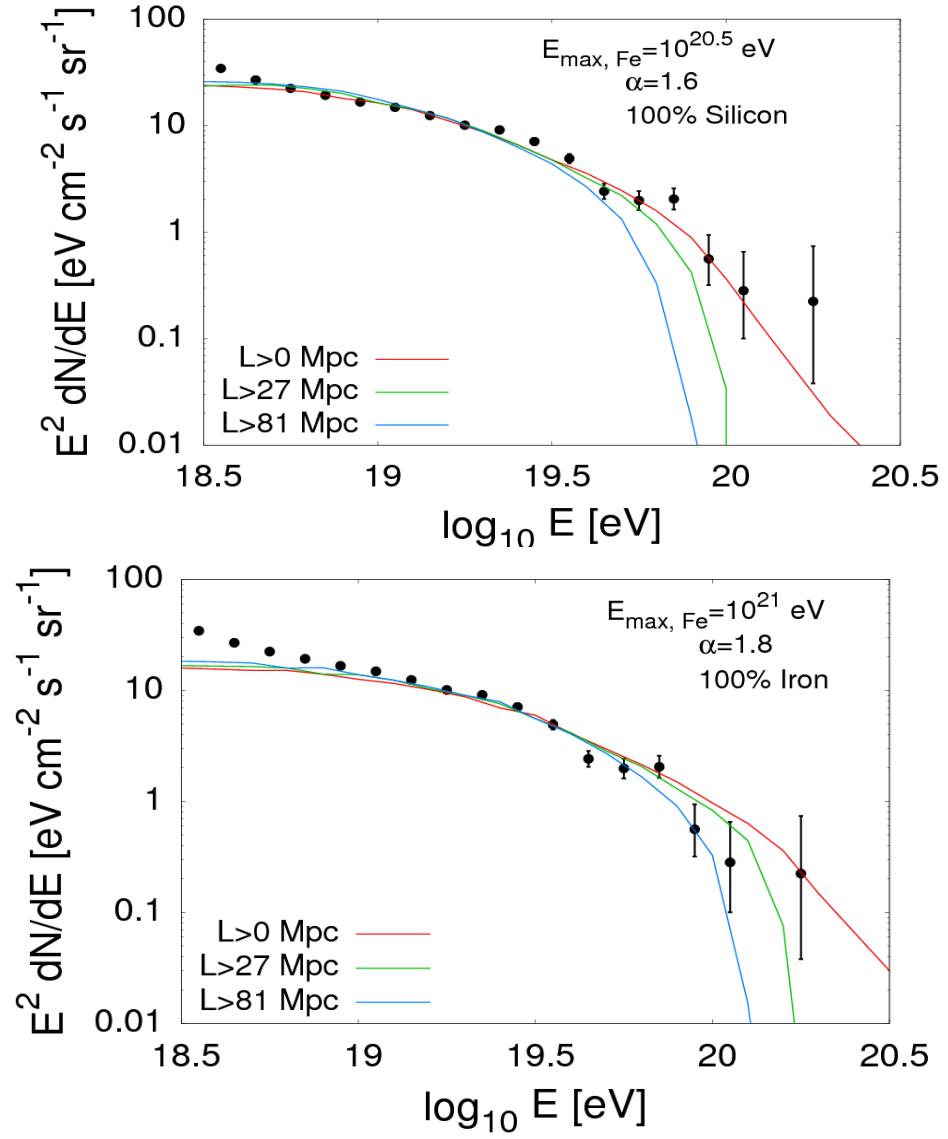
$$f_q(t) = \sum_{n=q}^m \frac{\tau_q \tau_n^{m-q-1}}{\prod_{p=q}^m (\tau_n - \tau_p)} e^{-\frac{t}{\tau_n}} f_n(0)$$



How Far is the Nearest Source?

Silicon- L<60 Mpc

Iron- L<80 Mpc



De Marco et al. (2006), 0603615

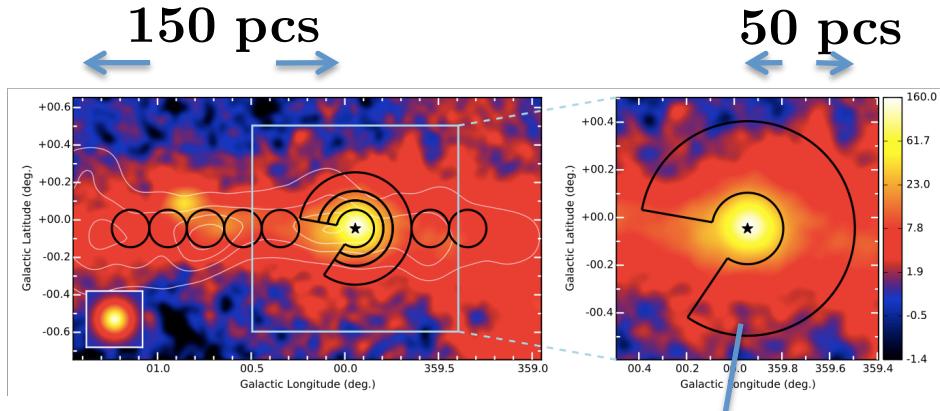
Taylor et al. (2011), 1107.2055

Fargion et al. (2015), 1412.1573

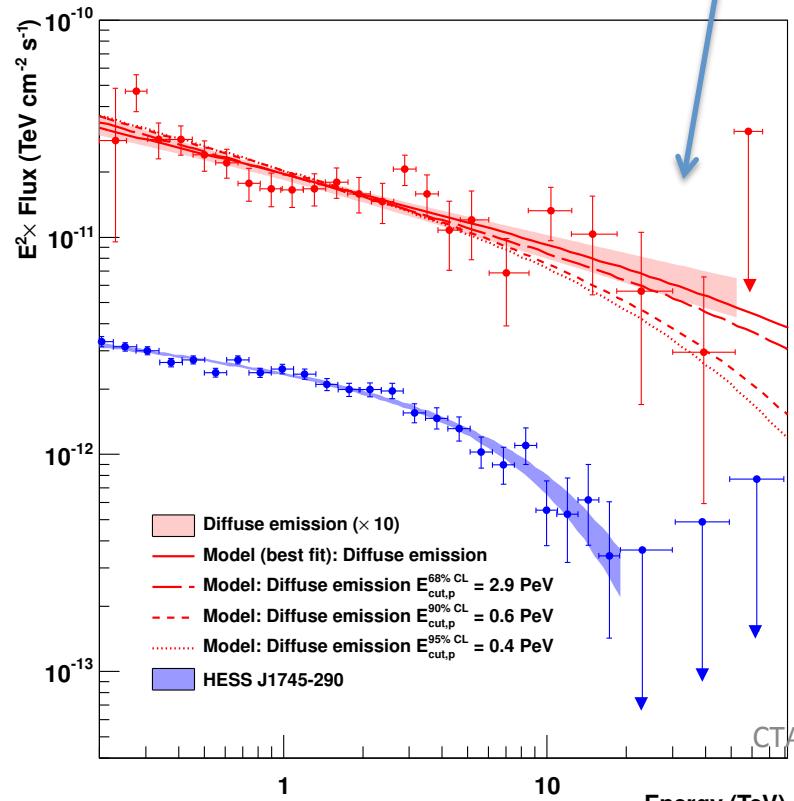
CTA Symposium- Andrew Taylor

Galactic Center- Sgr A*

HESS Coll. Nature 531 (2016) 476



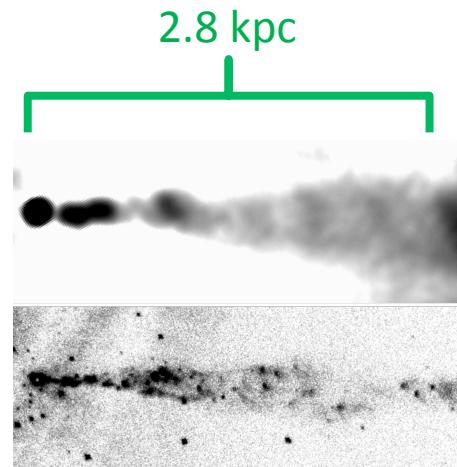
- Analysis of Sgr A* ‘point source’ at Galactic center
- Inflection evident in spectrum around 100 GeV revealing presence of new hard component



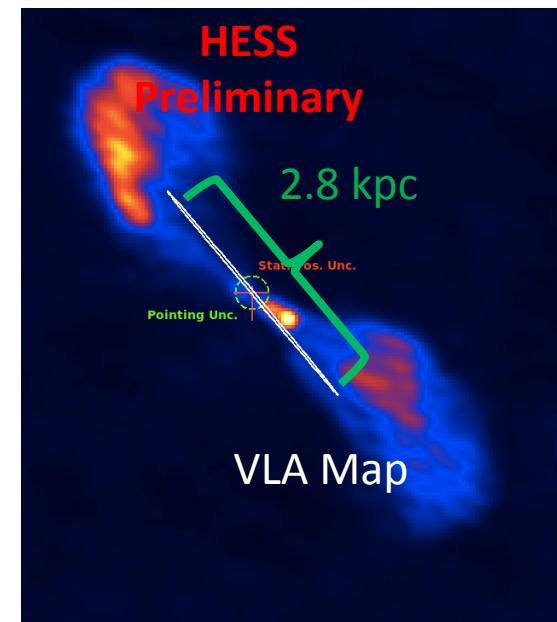
- Further Questions:
- 1) Maximum energy of cosmic rays produced? (info from Auger)
 - 2) Stability of the source Power?

Max Energy on Smaller Scales of Cen A?

Hardcastle et al. (1103.1744)



HESS Coll. TeVPA 2018



$$t_{\text{acc}} = \eta \frac{R_{\text{lar}}}{c \beta^2}$$

$$E_{\gamma}^{\text{sync}} \approx \frac{9}{4} \eta^{-1} \beta^2 \frac{m_e}{\alpha}$$

$$t_{\text{cool}} = \frac{9}{8\pi\alpha} \left(\frac{m_e}{E_{\gamma}^{\text{sync}}} \right) t_{\text{lar}}$$

For $\beta_{\text{scat.}} \approx 10^{-1}$, $\eta \approx 10^3$

Maximum energy
(Hillas criterion)

$$E_{\max} \approx 10^{15} \text{ eV}$$

Other (Local) Candidates- GRBs

New Information About sGRBs

Acceleration within remnant timescale constraint ($t_{\text{acc}} < 100$ days)

$B > 0.02 \text{ mG}$

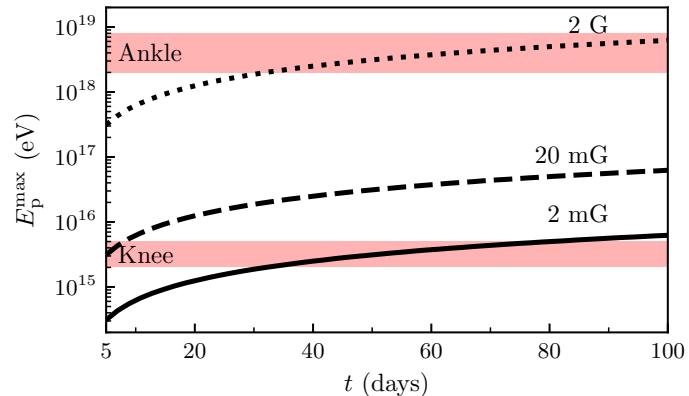
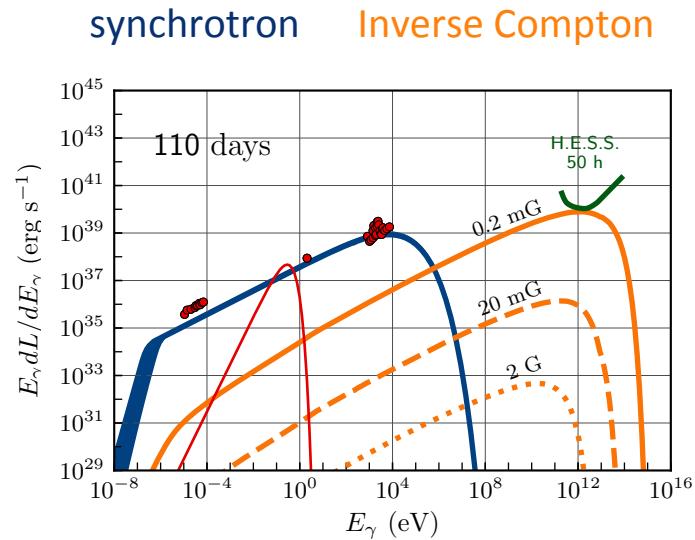
Absence of break in synchrotron emission spectrum up to X-ray energies constraint

$B < 2 \text{ mG}$

...alternatively, synchrotron emitting electrons may be always “fresh”, or the injection spectrum from the source may be very hard and the electrons observed cooled

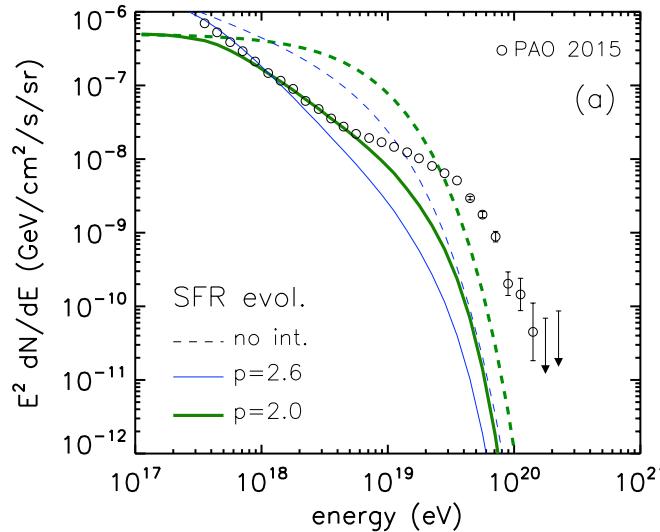
Kimura et al. 2018 (1807.03290)

Rogrigues et al. 2018 (1806.01624)

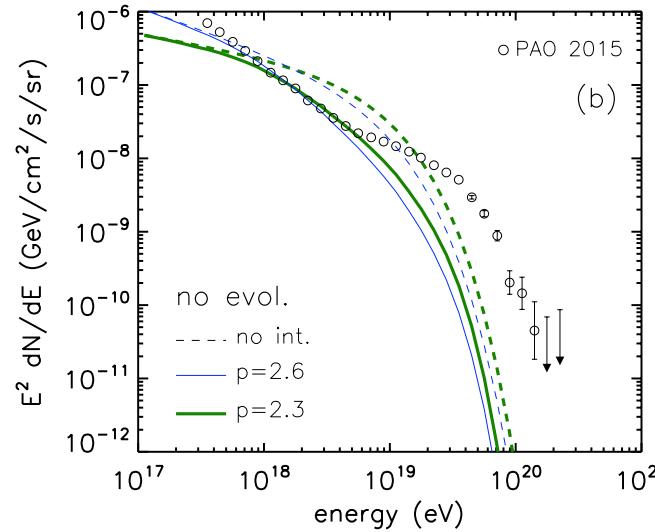


The Origin of Protons Below the Ankle

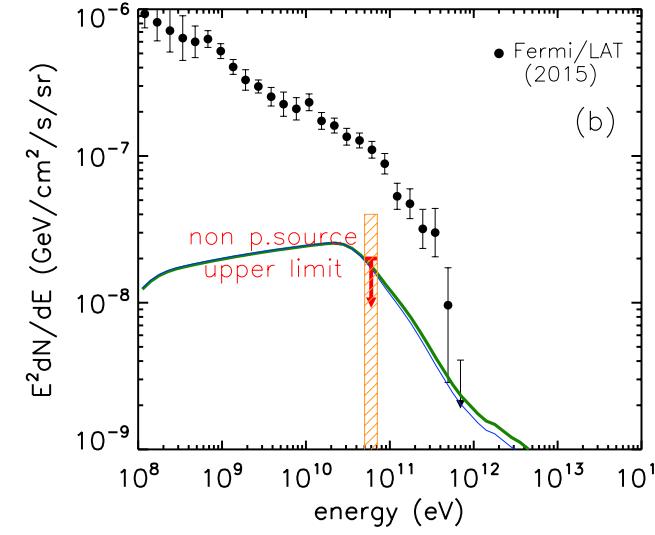
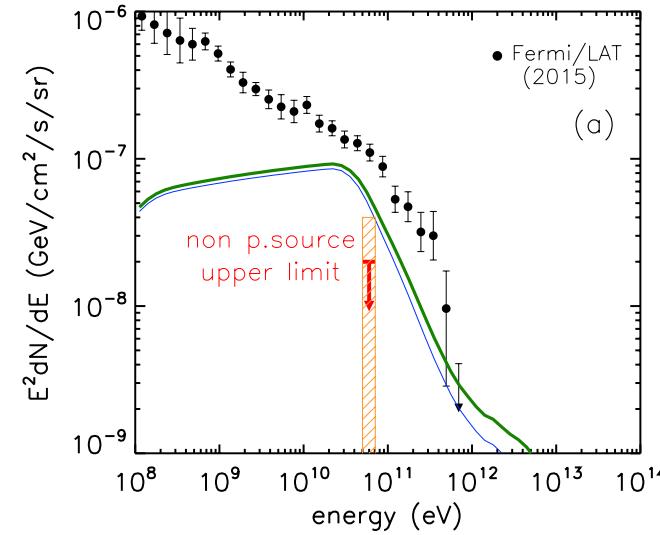
SFR evolution scenario



no evolution scenario



Note- IGRB contribution from cascade losses rather independent of source spectra



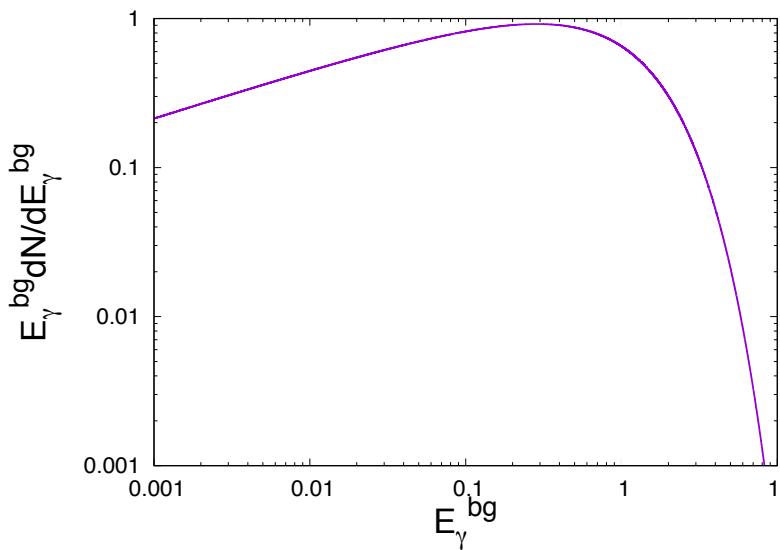
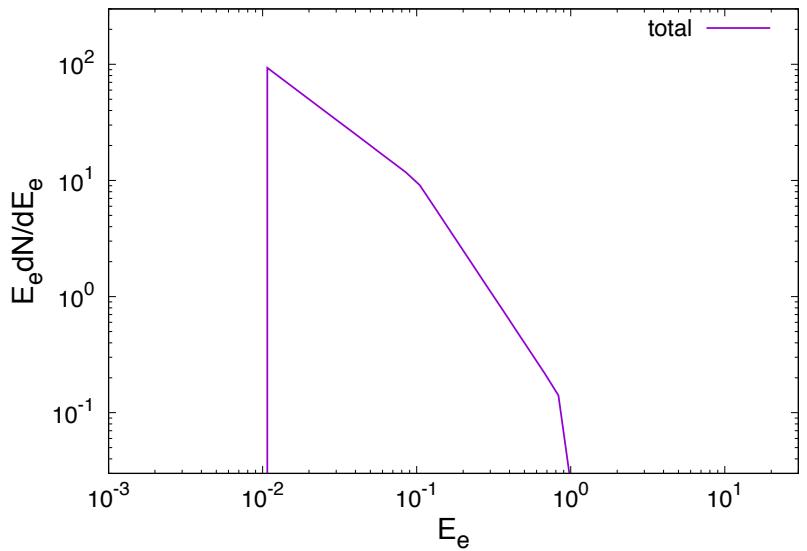
Liu et al. (2016), 1603.03223

Decerprit et al. (2011), 1107.3722

Gelmini et al. (2012), 1107.1672

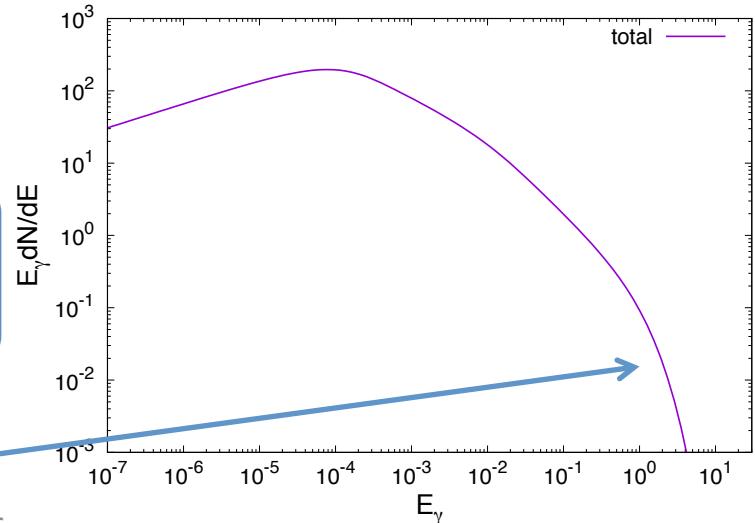
Crotron et al. 2019
(1801.10172)

Future Probes- Cutoff Region



$$\mathbf{E}_\gamma = \gamma_e^2 \left(\frac{\mathbf{B}}{\mathbf{B}_{crit}} \right) \mathbf{m}_e$$

$$E_\gamma \frac{dN}{dE_\gamma}_{tot} = \int \left(\frac{E_\gamma}{E_e^2} \right) \frac{dN}{dE_\gamma} \left(\frac{E_\gamma}{E_e^2} \right) E_e \frac{dN}{dE_e} dE_e$$



Possibility to probe cutoff region

CTA Symposium- Andrew Taylor

Cut-off Generation- A Simple Case

- Bohm diffusion ($q=1$) + only escape results in simple exponential cutoff.
- Some simplifications to the transport equation:

$$\cancel{\frac{\partial f}{\partial t}} = \nabla_p \cdot \left[(D_{pp} \nabla_p f) - \cancel{\frac{p}{\tau_{loss}(p)} f} \right] - \frac{f}{\tau_{esc}(p)} + \frac{Q}{p^2}$$

Steady state

No losses

Delta injection

$$\frac{\partial^2 f}{\partial p^2} + \frac{3}{p} \frac{\partial f}{\partial p} - \left(\frac{1}{D_0 \tau_0} \right) f = \delta(p)$$

Cutoff comes from balancing
1st and 3rd term

$$f \propto A e^{-p/p_\tau}$$

$$q = 1, r = 0, \rightarrow \beta_e = 1$$

Cut-off Shape- Emission Dependence

$$\frac{dN}{dE_e} \propto E_e^{-\Gamma} e^{-(E_e/E_{max})^{\beta_e}}$$

$$\frac{dN}{dE_\gamma} \propto E_\gamma^{-\Gamma} e^{-(E_\gamma/E_{max})^{\beta_\gamma}}$$

Recall generally, $\beta_e = 2 - q - r$

- Different emission processes dictate different relation between electrons and gamma rays

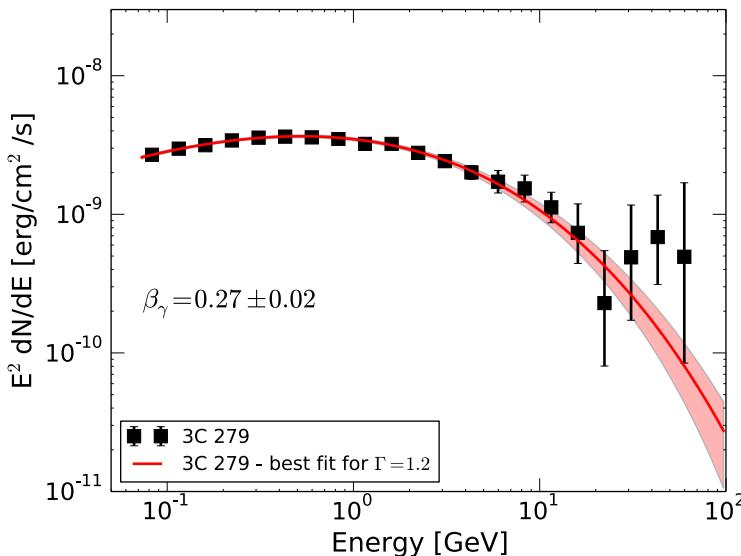
e.g.

- Synchrotron/IC Thomson: $\beta_\gamma = \frac{\beta_e}{\beta_e + 2}$
- SSC: $\beta_\gamma = \frac{\beta_e}{\beta_e + 4}$
- IC (Klein Nishina) $\beta_\gamma = \beta_e$

Good measurement of gamma ray cut-off can give insight into the acceleration environment

Observation of Cut-offs in Gamma-ray Spectra

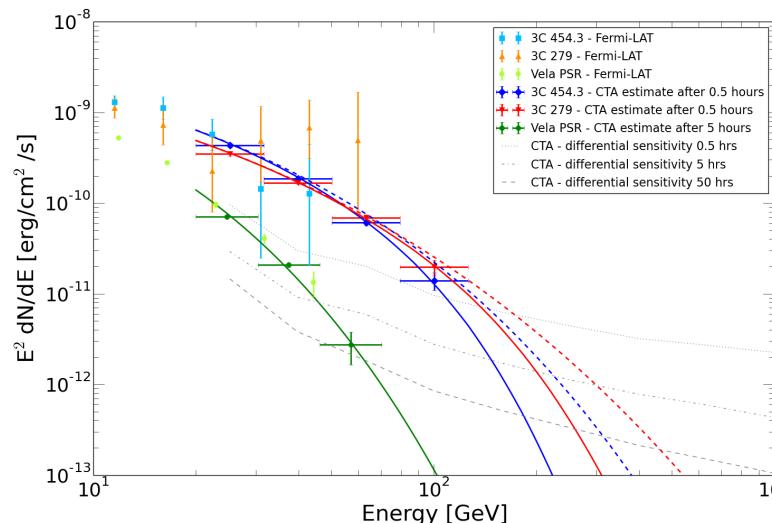
- 2nd Brightest AGN Flare-
3C 279 June 2015



Romoli et al.,
**Astropart.Phys. 88 38-45
(2017)**

| Parameter | $\Gamma = 1.2$ |
|---------------------------------|-------------------------------|
| N [ph/cm ² /s/GeV] | $(2.8^{+0.8}_{-0.6}) 10^{-4}$ |
| Γ (fixed) | 1.2 |
| E_c [GeV] | $(8.4^{+6.6}_{-4.1}) 10^{-3}$ |
| β_γ | 0.27 ± 0.02 |
| E_s (fixed) [GeV] | |

Values obtained on a 3 days integration
Note- X-ray observations during flare
indicated that $\Gamma = 1.17 \pm 0.06$



Study using the expected
CTA performance
Fermi data integrated
over 3 days
Constraint on
parameter at 10% level
obtained during only 0.5
hr flare!

Conclusion

- Nearby extragalactic cosmic ray sources must exist
- Further probe of the Galactic center still needed to understand nature and limitations of source
- The probe of nearby candidates is helping to put the pieces in place about how these accelerators operate
- The time domain holds key potential for probing sources, which CTA is particularly suited to take advantage of

A Simple Case (II)- q=1, only escape

- Rearranging the terms (and explicitly stating the dependences from p of the parameters):

$$\frac{1}{p^2} \frac{\partial}{\partial p} \left(p^2 D_0 \frac{p}{p_0} \frac{\partial f}{\partial p} \right) - \frac{f}{\tau_{\text{esc}}(p)} = \delta(p), \quad \tau_{\text{esc}}(p) \propto p^{-1}$$

$$\frac{\partial^2 f}{\partial p^2} + \frac{3}{p} \frac{\partial f}{\partial p} - \left(\frac{1}{D_0 \tau_0} \right) f = \delta(p)$$

Cutoff comes from balancing
1st and 3rd term

$$f \propto A e^{-p/p_\tau}$$

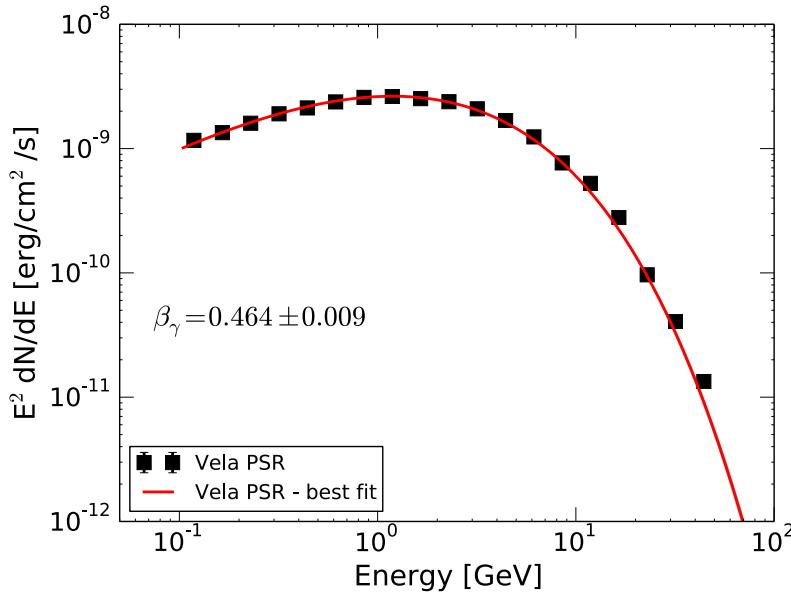
Recall generally, $\beta_e = 2 - q - r$

$$q = 1, \quad r = 0, \quad \rightarrow \quad \beta_e = 1$$

(Note- energy losses for the $r = 0$ case will not alter this result)

Observation of Cut-offs in Gamma-ray Spectra

- Test case- Vela Pulsar (brightest source)



$$\frac{dN}{dE_\gamma} \propto E_\gamma^{-\Gamma} e^{-(E_\gamma/E_{\max})^{\beta_\gamma}}$$

| Parameter | Value |
|---------------------------------|----------------------------------|
| N [ph/cm ² /s/GeV] | $(1.39_{-0.10}^{+0.12}) 10^{-5}$ |
| Γ | 1.019 ± 0.011 |
| E_c [GeV] | 0.238 ± 0.016 |
| β_γ | 0.464 ± 0.009 |
| E_s (fixed) [GeV] | 0.83255 |

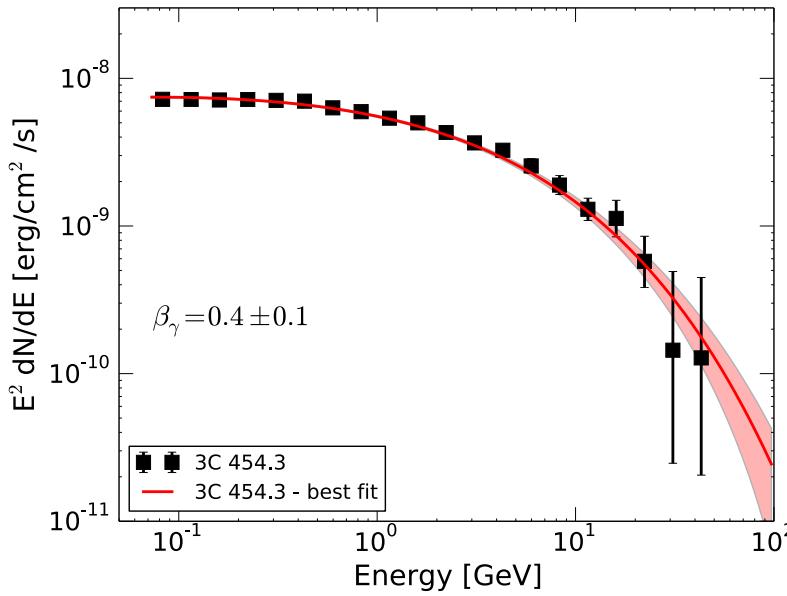
- Note- MCMC method used to explore ‘good-fit’ region. This has the benefit of being stable on the landscape being explored

Romoli et al., **Astropart.Phys.** **88** 38-45
(2017)



Observation of Cut-offs in Gamma-ray Spectra

- Brightest AGN Flare-
3C 454 Nov 2010



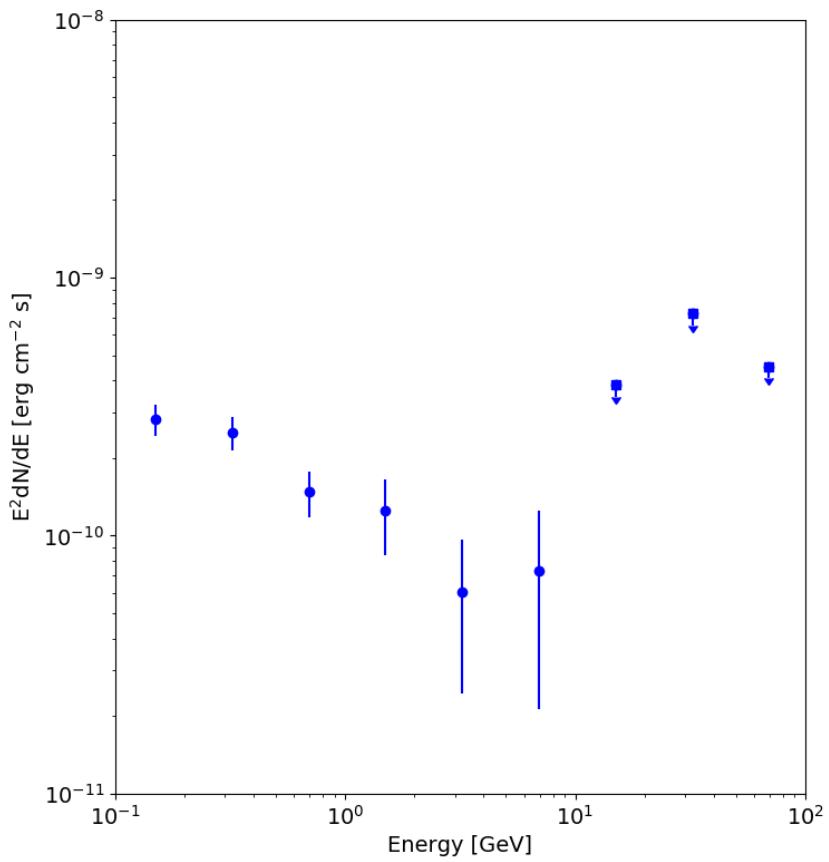
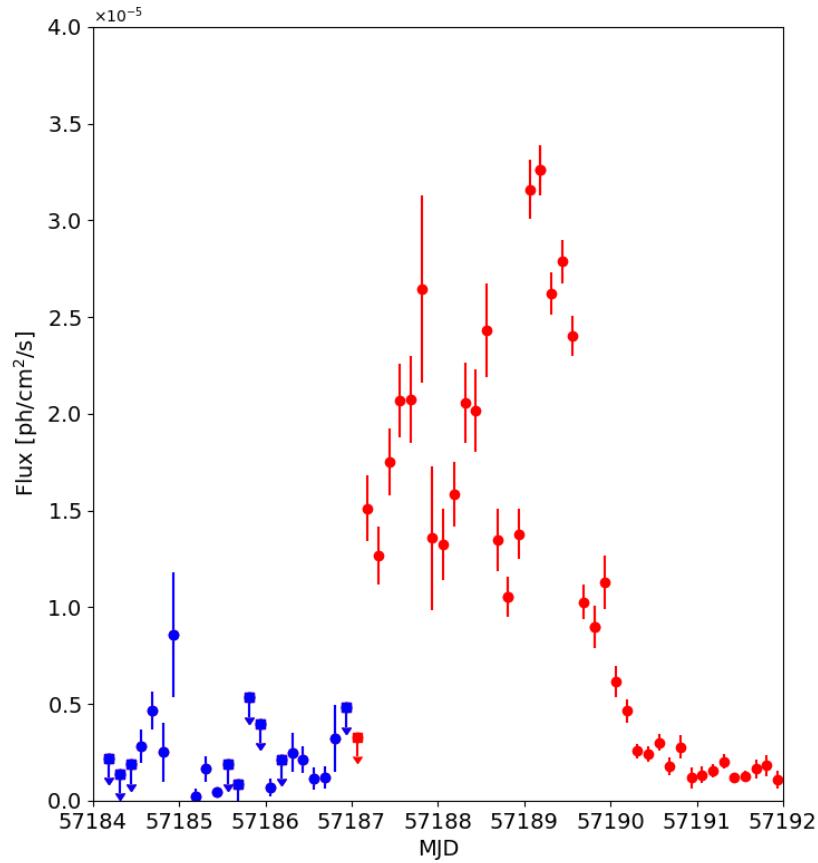
Romoli et al., **Astropart.Phys.** **88** 38-45
(2017)



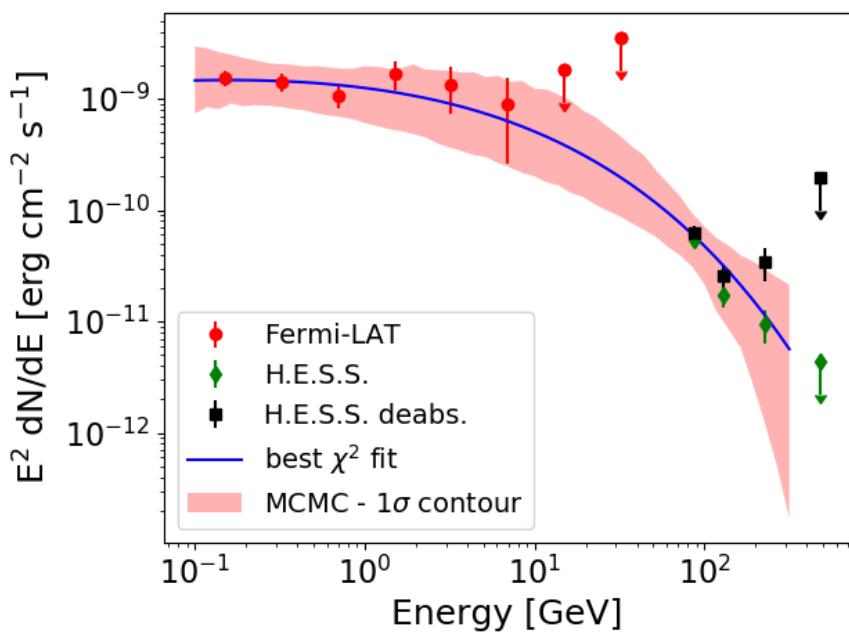
| Parameter | Value |
|---------------------------------|-------------------------------|
| N [ph/cm ² /s/GeV] | $(4.7^{+3.9}_{-1.2}) 10^{-5}$ |
| Γ | $1.87^{+0.08}_{-0.12}$ |
| E_c [GeV] | $1.1^{+1.6}_{-0.9}$ |
| β_γ | 0.4 ± 0.1 |
| E_s (fixed) [GeV] | 0.41275 |

- Indicating a cut-off value of the primary particles around 1 GeV
- Caveats:**
 - Values obtained on a 7 days integration (for statistics)
 - Spectrum variable during the flare -> superposition effects?

3C 279 June 2015 Flare- Temporal Evolution



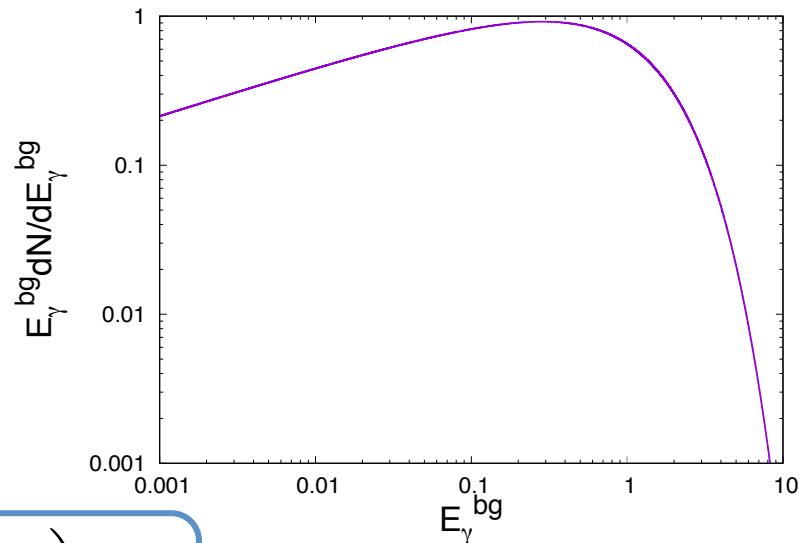
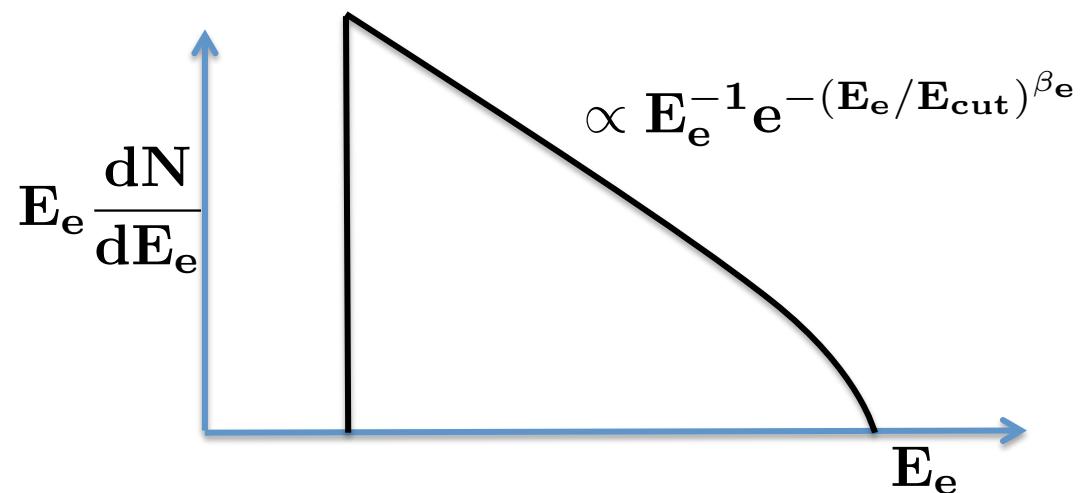
Can We Do Better Already? Fermi + H.E.S.S.II Fit



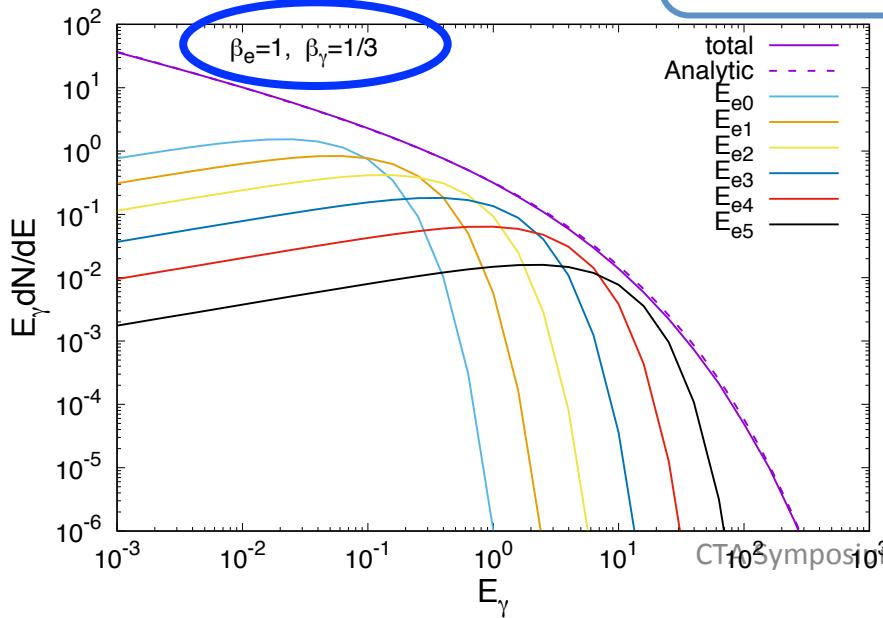
| Parameter | MCMC fit |
|---|--|
| $\log_{10} N_0 [\text{ph}/\text{cm}^2/\text{s}/\text{GeV}]$ | $(-4.75^{+0.91}_{-0.24}) \times 10^{-5}$ |
| Γ | $(1.93^{+0.29}_{-0.41})$ |
| $\log_{10} E_c [\text{GeV}]$ | $0.13^{+1.33}_{-2.82}$ |
| β_γ | $0.34^{+0.32}_{-0.14}$ |

- Joint fit of Fermi-LAT data (9 hours centred on HESSII obs.) taken on night 2
- $\beta_\gamma = 0.34^{+0.32}_{-0.14}$

Cut-off Shape- Electrons & Photons

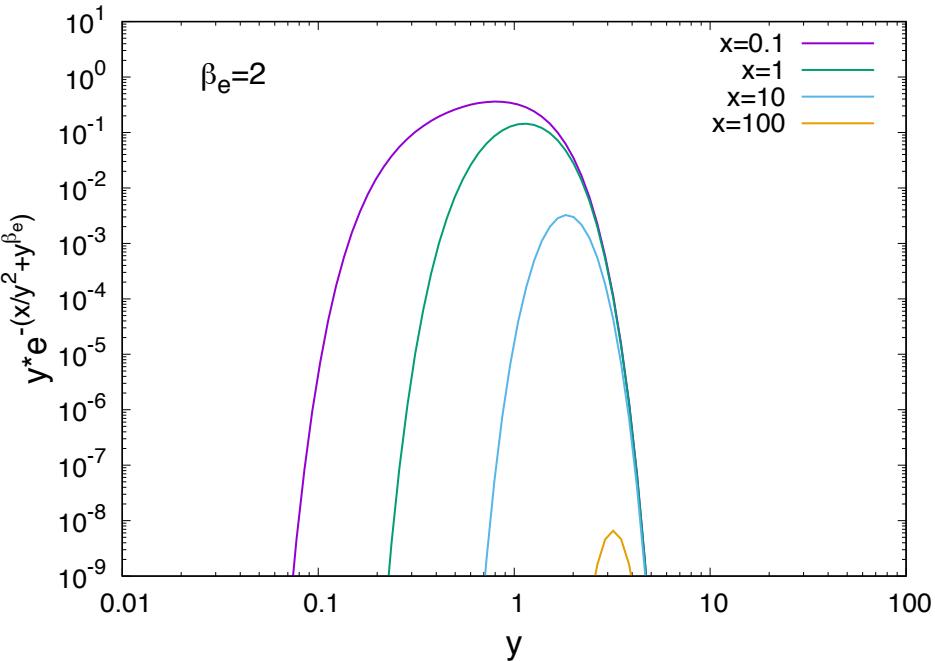
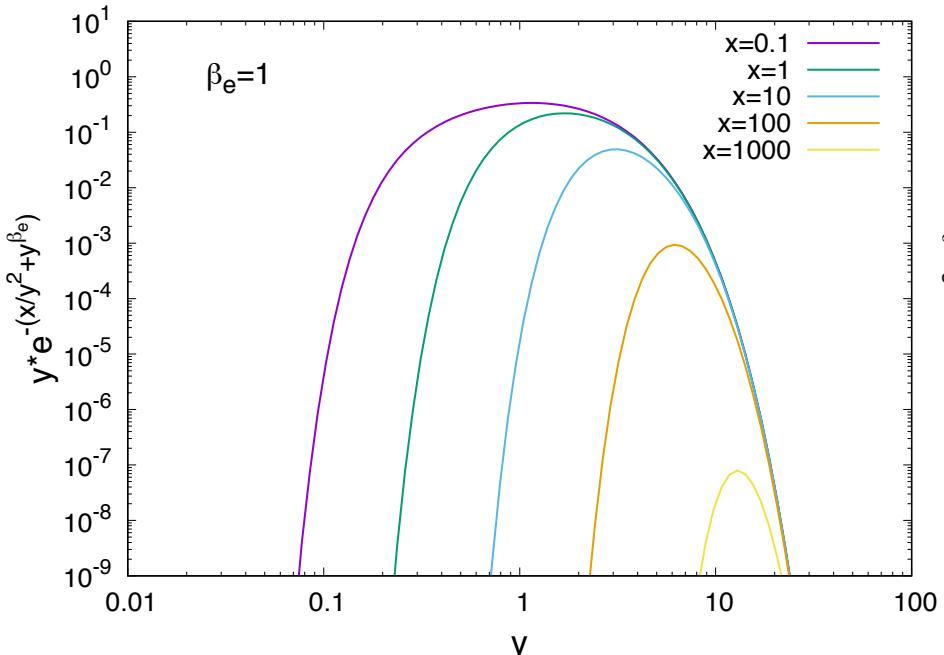


$$E_\gamma = \gamma_e^2 \left(\frac{B}{B_{crit}} \right) m_e$$



$$f(x) = \int_0^\infty e^{-(x/y^2)} e^{-y^{\beta_e}} dy$$

Integrand-



$$y^2 \left(y^{\beta_e} - \frac{1}{\beta_e} \right) = \frac{2x}{\beta_e}$$

$$y^2 \approx \left(\frac{2x}{\beta_e} \right)^{\frac{2}{\beta_e + 2}}$$

$$\frac{x}{y^2} \approx x^{\frac{\beta_e}{\beta_e + 2}}$$

$$\beta_\gamma = \frac{\beta_e}{\beta_e + 2}$$

CTA Symposium- Andrew Taylor

Further Acceleration Further Out?

$$R_{\text{lar}} = \frac{\beta_{\text{scat}}}{\eta} R$$

→ $R_{\text{Lar}}(10^{18} \text{ eV p}) \approx 100 \text{ pc}$

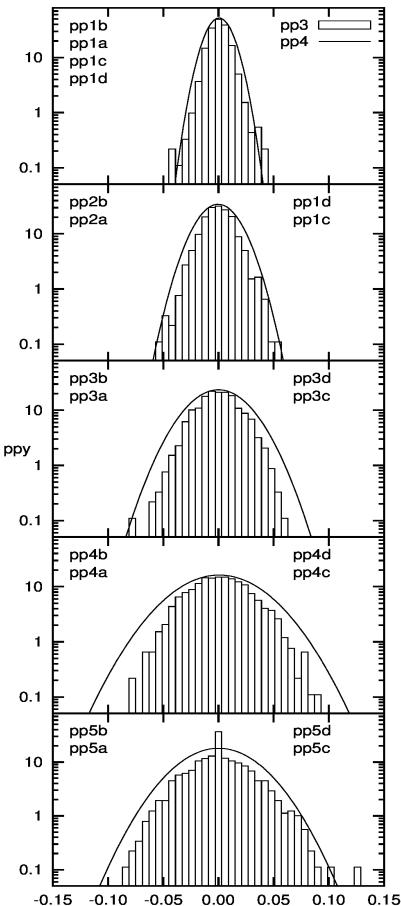
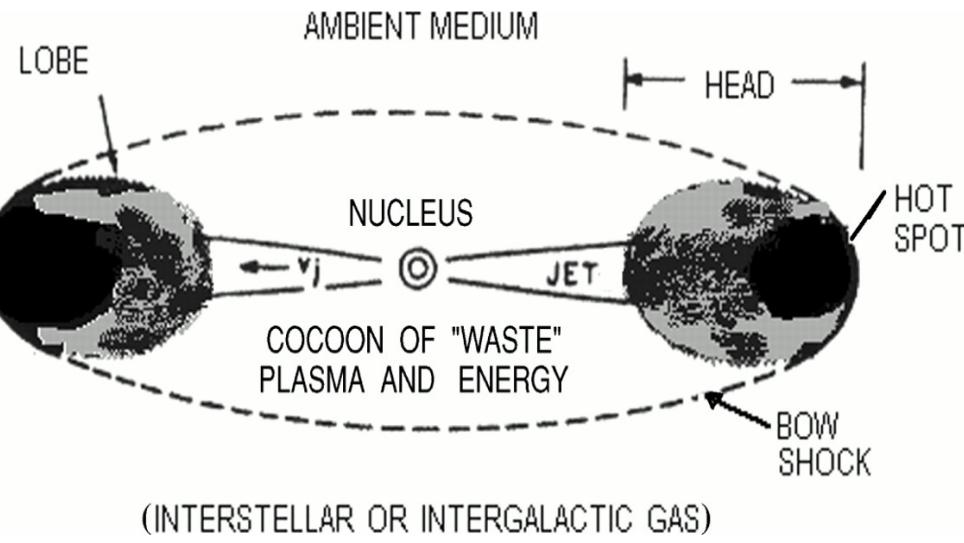


Diagram taken from Ferrari -1998



(INTERSTELLAR OR INTERGALACTIC GAS)

$$E_{\text{max}} \approx 10^{18} \text{ eV}$$

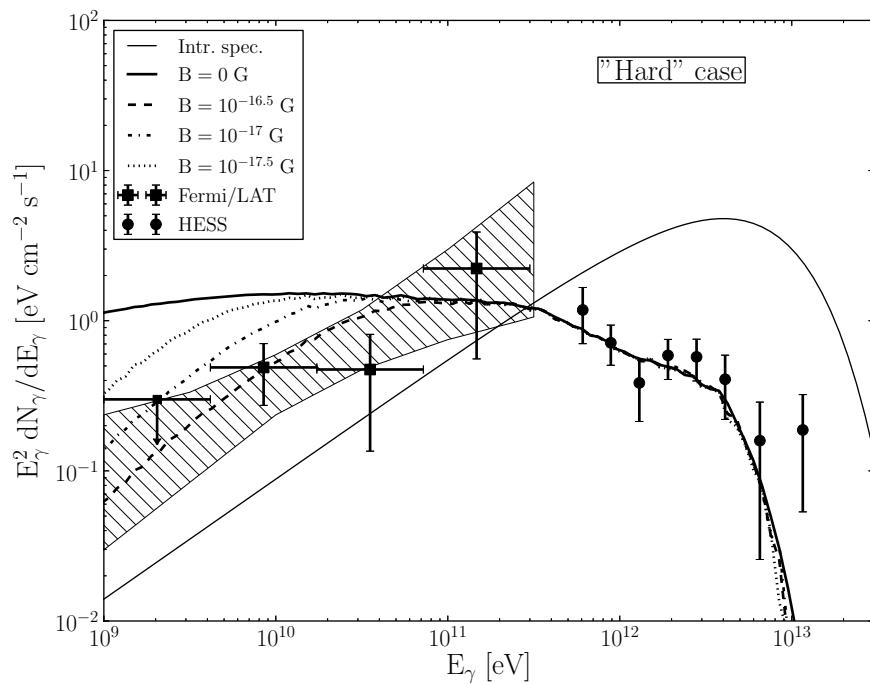
$$t_{\text{acc}} \approx 0.1 \text{ Myr}$$

O'Sullivan et al. 2009 (0903.1259)

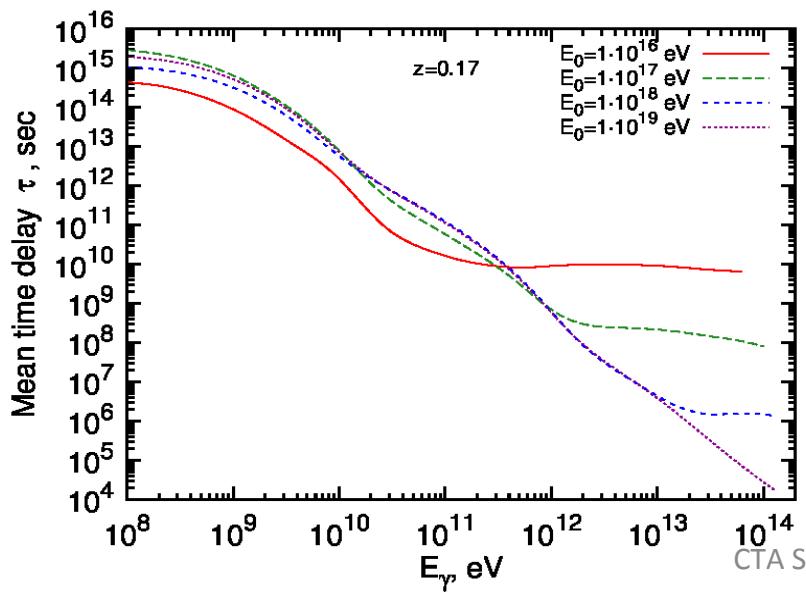
Matthews et al. 2019 (1902.10382)

Future Probes- Temporal Structure

Possibility that emission comes from much higher energy emission (potentially from proton losses.....)



Prosekin et al. (1203.3787)



CTA Symposium- Andrew Taylor

Taylor et al. (1101.0932)

